



GDR
H y d r o G E M M

Hydrogen storage in salt caverns: evolution of salt permeability and hydromechanical modelling at the cavern scale under cycling conditions.

Dennys Coarita, **Fabrice Golfier**, Dragan Grgic, Farah Al Sahyouni, Mountaka Souley, Long Cheng

Université de Lorraine, CNRS, UMR 7359 GeoRessources, Nancy (France)

Institut National de l'Environnement Industriel et des Risques (Ineris), Nancy, France

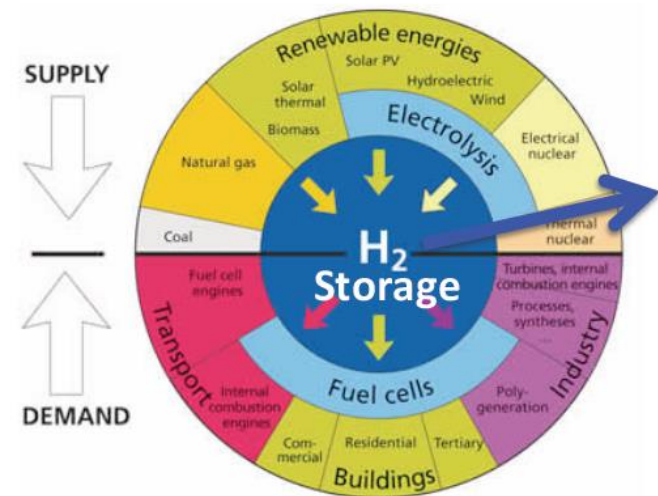
Funding : ROSTOCK-H Geodénergies project, Grand Est Region Council, Carnot ICEEL



GENERAL CONTEXT: energy transition and H₂ storage

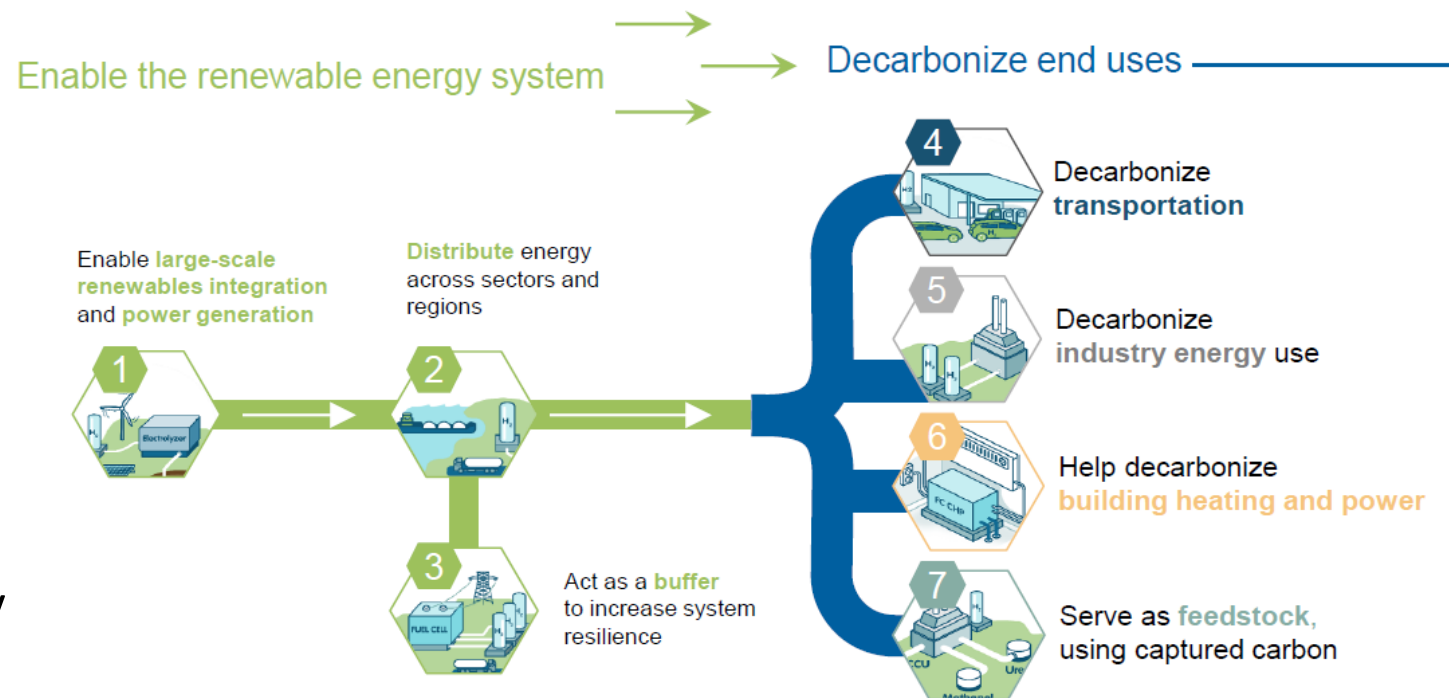
Hydrogen: versatile and promising carrier for energy transition

Currently produced from fossil fuels but from renewable energies in the future (green hydrogen)



Clean hydrogen (green and/or blue) should play a key role in the world transition to achieve carbon neutrality before 2050

- ✓ large-scale electricity storage
- ✓ decarbonize uses that are difficult to electrify
- ✓ use in industrial processes



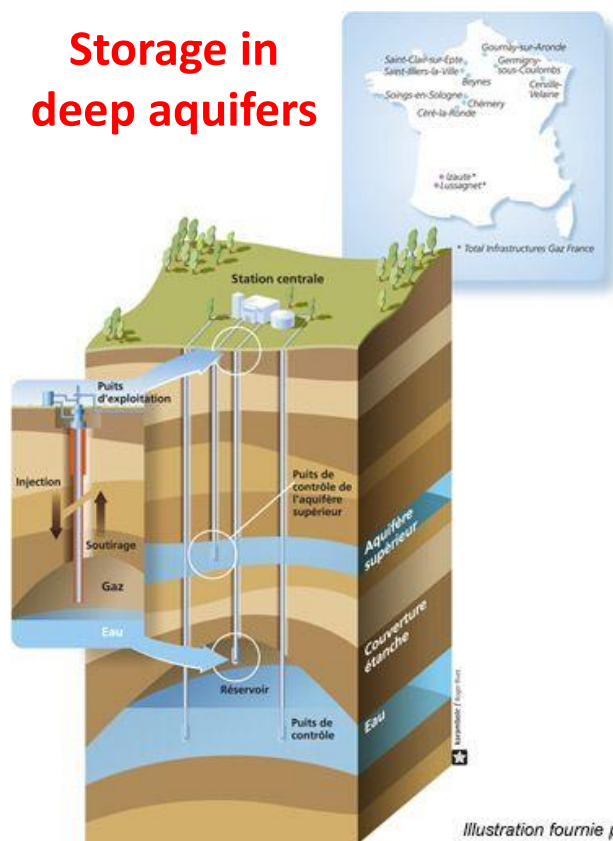
GENERAL CONTEXT: massive H₂ storage

Storing hydrogen in large quantities is the best option

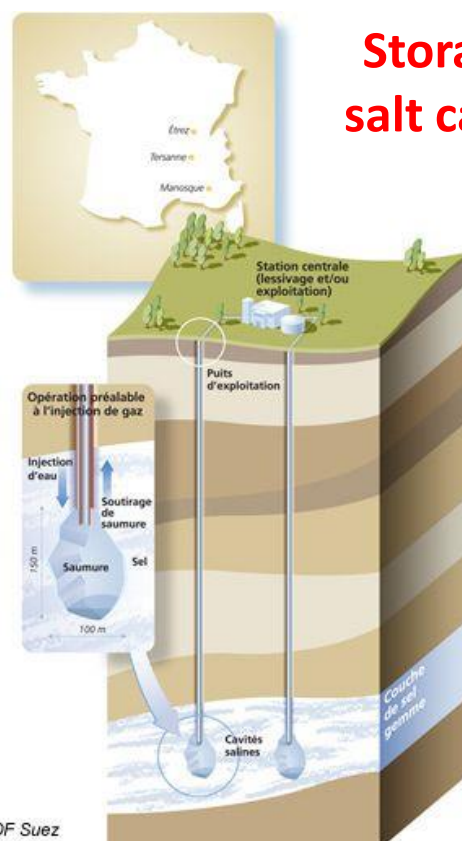
Massive Storage ➡ **Underground storage**

2 solutions

Storage in deep aquifers



Storage in salt caverns



Storage in salt cavern

Most secure and economical solution for large volumes of H₂

Only 4 salt caverns in operation worldwide (e.g., Spindletop USA, Teeside UK...)

but many industrial pilot sites in preparation worldwide (e.g., Etrez, France)

PROBLEM

Specific constraints compared to conventional gas storage

- High danger and mobility of H_2 (low density and viscosity, high diffusion)
- Need to accurately predict the extent of potential leaks of H_2 to ensure safe storage: **sealing of salt?**

Strong temperature (0-50 °C) and gas pressure (2-6 MPa for a 350 m deep cavern) variations in the salt cavern depending on the injection-production cycles (depends on usage)

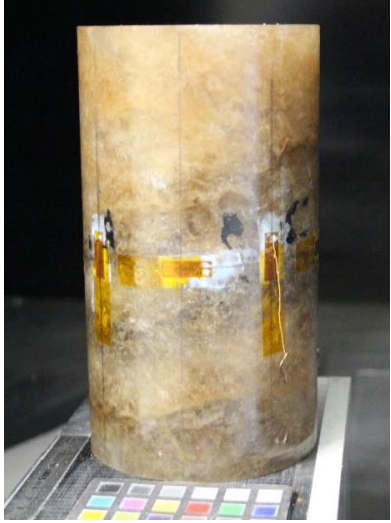
→ Thermo-mechanical damage of rock salt in the near field (close to the wall) → permeability ↗

FOCUS OF THE STUDY

Characterize the ThermoHydroMechanical behaviour of rock salt

- Mechanical properties of initial material
- Impact of damage processes on rock salt permeability
- Viscoplasticity effect (self-healing process)
- Impact of mechanical (static and dynamic) and thermal (dynamic) fatigue

MATERIALS AND METHODS



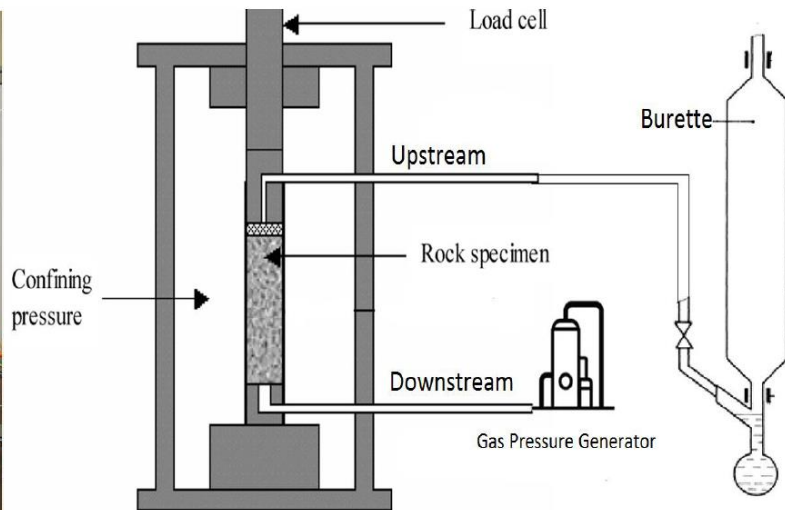
Cylindrical sample of rock salt
($\phi = 100$ mm; $H = 200$ mm)

Rock salt samples

- Salt bed of the Alsace potash mines (530 m depth) in the East region of France (Stocamine site for ultimate waste storage)
- Sannoisian-Oligocene (Cenozoic) geological stage
- Considered as a natural analogue of salt caverns used for in-situ H_2 storage
- Very low initial porosity ($\sim 1\%$) composed mainly of intrapores (nanometric size) that connect dispersed cracks and macropores

Experimental device

- Large scale triaxial compression cell
- Continuous measurement of deformations (strain gages)
- Continuous measurement of gas (He) permeability with the steady-state method
- Pressures controlled with precision syringe pumps



Experimental device

EXPERIMENTAL PROGRAM

Characterisation of the mechanical behaviour

- Hydrostatic compression tests
 - Uniaxial and triaxial compression tests
- Characterization of short-term deformation mechanisms (plasticity, micro-cracking damage)

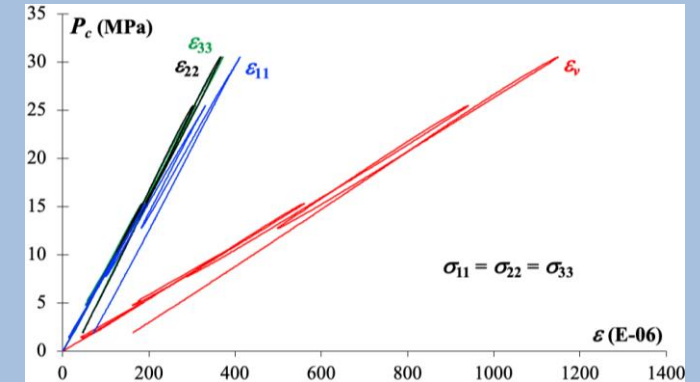
Gas permeability measurements

- Under hydrostatic loading and triaxial loading with different confining pressures
- Analyse the impact of the deformation mechanisms on permeability
- At different gas pressures
- Analyse the slippage (Klinkenberg) effect and obtain the intrinsic gas permeability
- Under dynamic (cyclic) and static (creep test) triaxial loading, and under dynamic thermal loading
- Analyse the impact of mechanical and thermal fatigue

MECHANICAL BEHAVIOUR OF ROCK SALT

Under hydrostatic loading

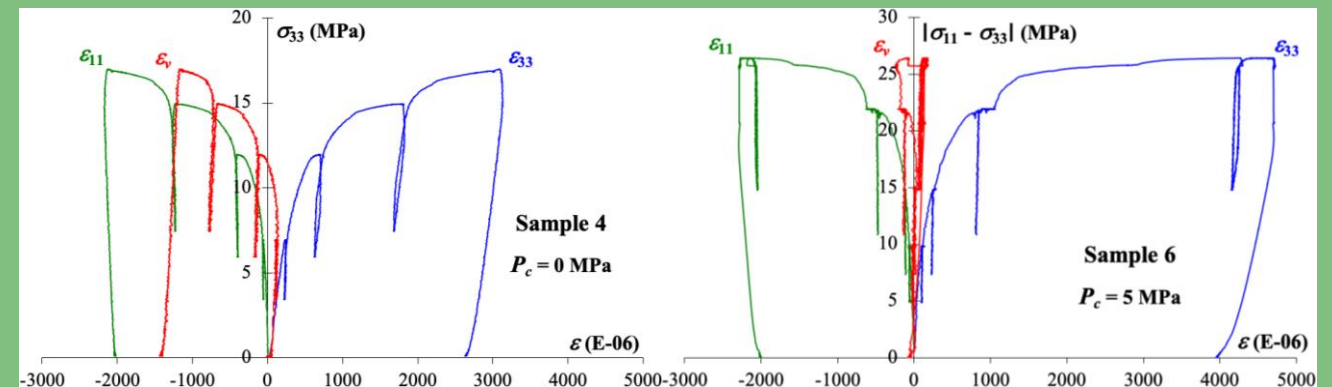
- **Isotropic** mechanical behaviour
 - The behaviour is **elastic** (irreversible) and linear up to 30 MPa
- No large initial cracks
- No poromechanical coupling (Biot's coefficient ~ 0)



Stress-strain curves of an hydrostatic compression test

Under deviatoric loading

- **Very low elastic limit** (yield strength)
- **Crack initiation and dilatancy thresholds increase** with the confining pressure (P_c)
- **Dilatant and irreversible microcracking damage** develops at low P_c (0 and 1 MPa)
- Behaviour becomes **fully plastic** (ductile) at high P_c (5 MPa)
- Dilatancy threshold $\leftrightarrow$ important increase of permeability due to cumulated microcracking damage



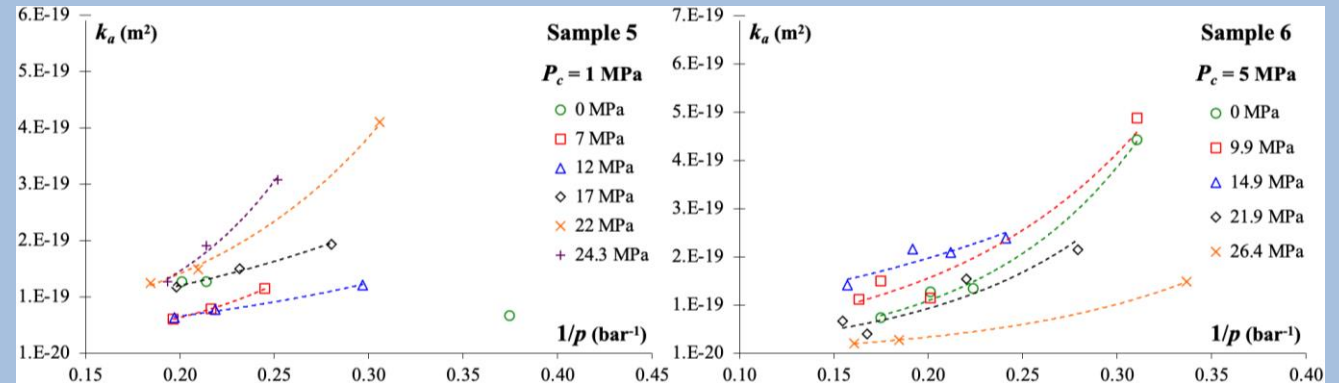
Stress-strain curves of of uniaxial (left) and triaxial (right) compression tests

INTRINSIC PERMEABILITY OF ROCK SALT

Klinkenberg (slippage) effect

- **Klinkenberg effect** (i.e., decrease of permeability with the increase in gas pressure) only observed for the less permeable (and then less initially damaged) samples
- When it appears, the gas flow falls in **transitional regime** in weakly damaged rock salt

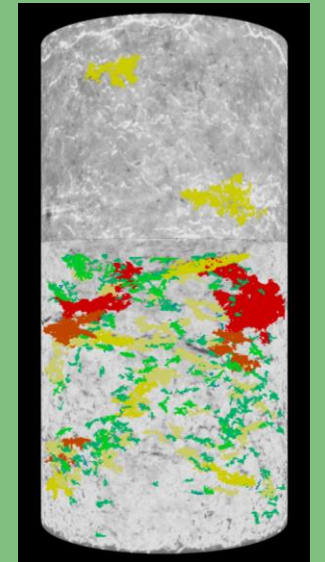
Evolution of apparent permeability k_a as a function of mean gas pressure at different deviatoric stress levels during triaxial compression tests



Initial permeability

- The initial intrinsic permeability of the studied rock salt ranges over more than 4 orders of magnitude: 10^{-16} to 5×10^{-21} m²
- Wide permeability range of the as-received samples due to the presence of cracks caused by the stress relaxation (induced by core drilling or cavity excavation) and sample preparation

Voids (cracks) distribution in a rock salt sample from X-ray tomography

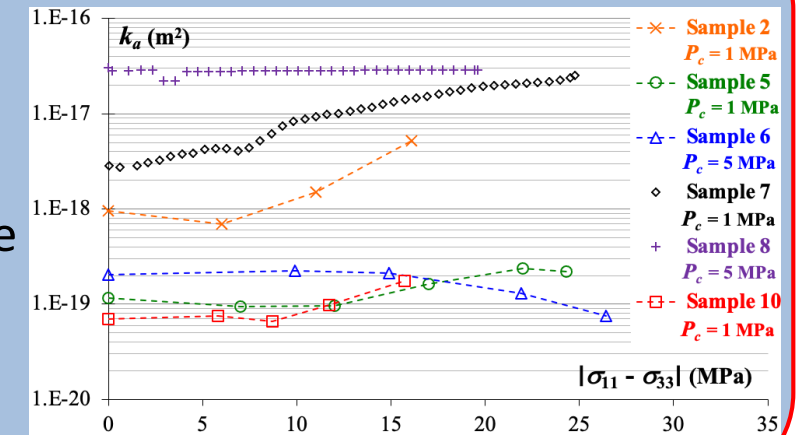


EVOLUTION OF APPARENT GAS PERMEABILITY WITH STRESS INCREASE

Under deviatoric loading

- At low confining pressure (1 MPa): small increase in gas permeability from the dilatancy threshold due to **microcracking damage**
- At high confining pressure (5 MPa): no increase in permeability because the material becomes **fully plastic** which practically eliminates microcracking and thus dilatancy

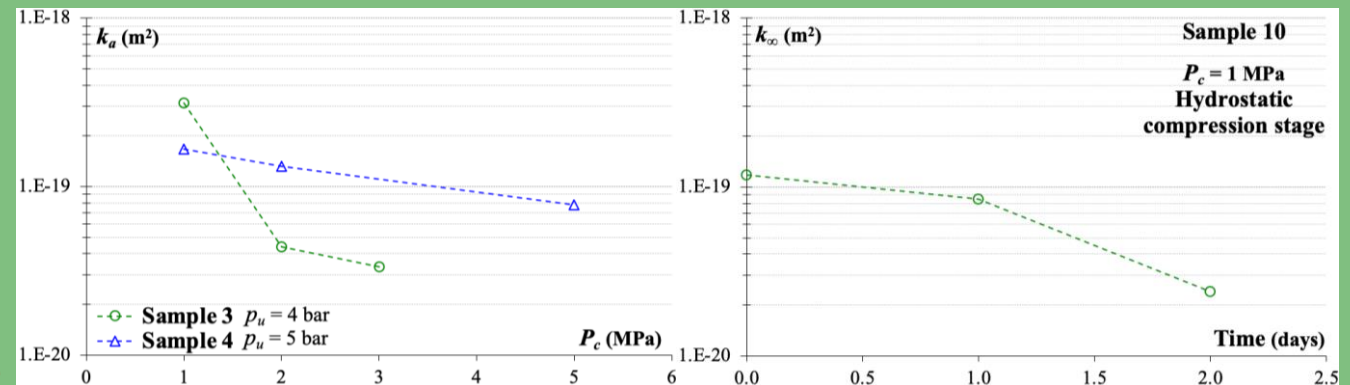
Evolution of apparent permeability k_a as a function of deviatoric stress



Under hydrostatic loading

- Gas permeability decreases because of the **closure of voids** (cracks and pores)
 - Decrease is irreversible if time of run is high enough
 - Due to the **self-healing process** (irreversible closure of cracks)
- Restoration of the permeability of undisturbed natural rock salt

Evolution of apparent permeability k_a as a function of confining pressure P_c and time

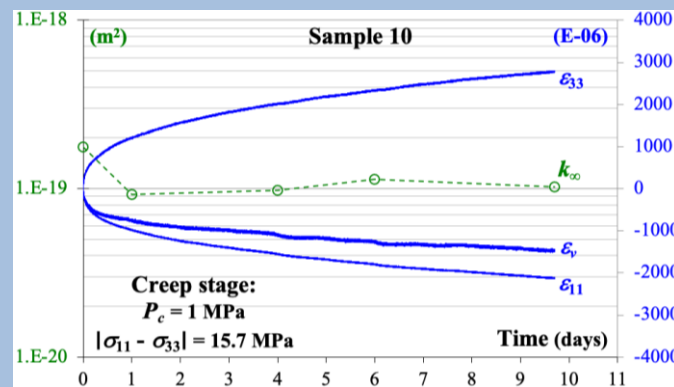


IMPACT OF MECHANICAL AND THERMAL FATIGUE ON ROCK SALT PERMEABILITY

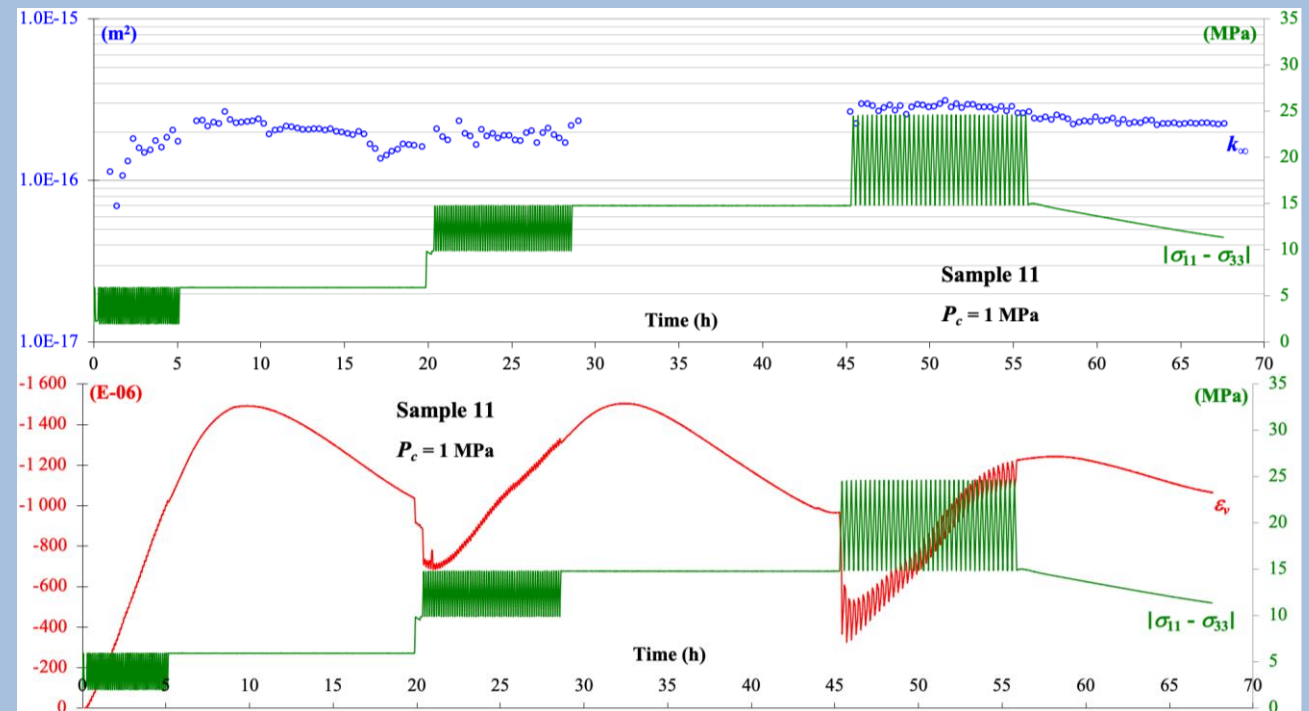
Static (creep test) and dynamic (cyclic) mechanical fatigue

- Volumetric dilatancy (microcracking damage) develops and increases slightly the permeability during dynamic fatigue
- Self-recovery reduces damage and decreases slightly the permeability during static fatigue (creep)

→ Different mechanisms involved in rock salt deformation during dynamic and static fatigue act in a competitive way to **annihilate any significant permeability evolution**



Evolution of permeability k and deformations as a function of time during a creep test

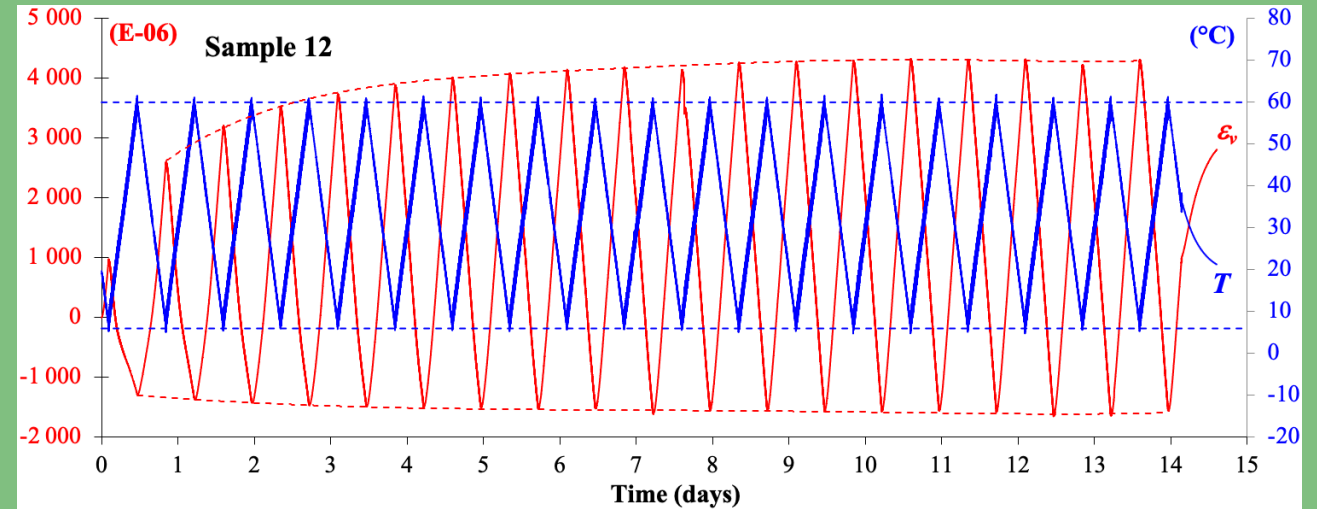


Evolution of permeability k and volumetric deformation ϵ_v as a function of time during the cyclic triaxial compression test

IMPACT OF MECHANICAL AND THERMAL FATIGUE ON ROCK SALT PERMEABILITY

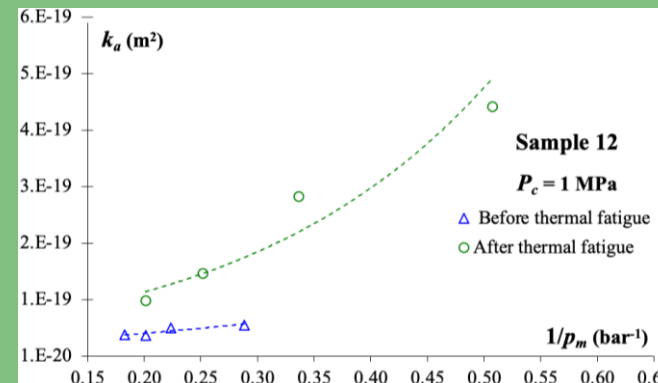
Cyclic thermal fatigue

- Small permeability increase due to the **microcracking damage** that develops at the **microscopic scale**
- Due to the **anisotropy of the thermal deformation** of rock forming minerals and to the **polycrystalline nature** of rock salt
→ Deformation heterogeneities and then differential thermal stresses and **microcracking damage**



Evolution of volumetric deformation ϵ_v as a function of time during the thermal cyclic fatigue test

Evolution of apparent permeability k_a as a function of mean gas pressure, before and after the thermal cyclic fatigue test



Intrinsic permeability
Before test: $2.7 \times 10^{-20} m^2$
After test: $4.9 \times 10^{-20} m^2$

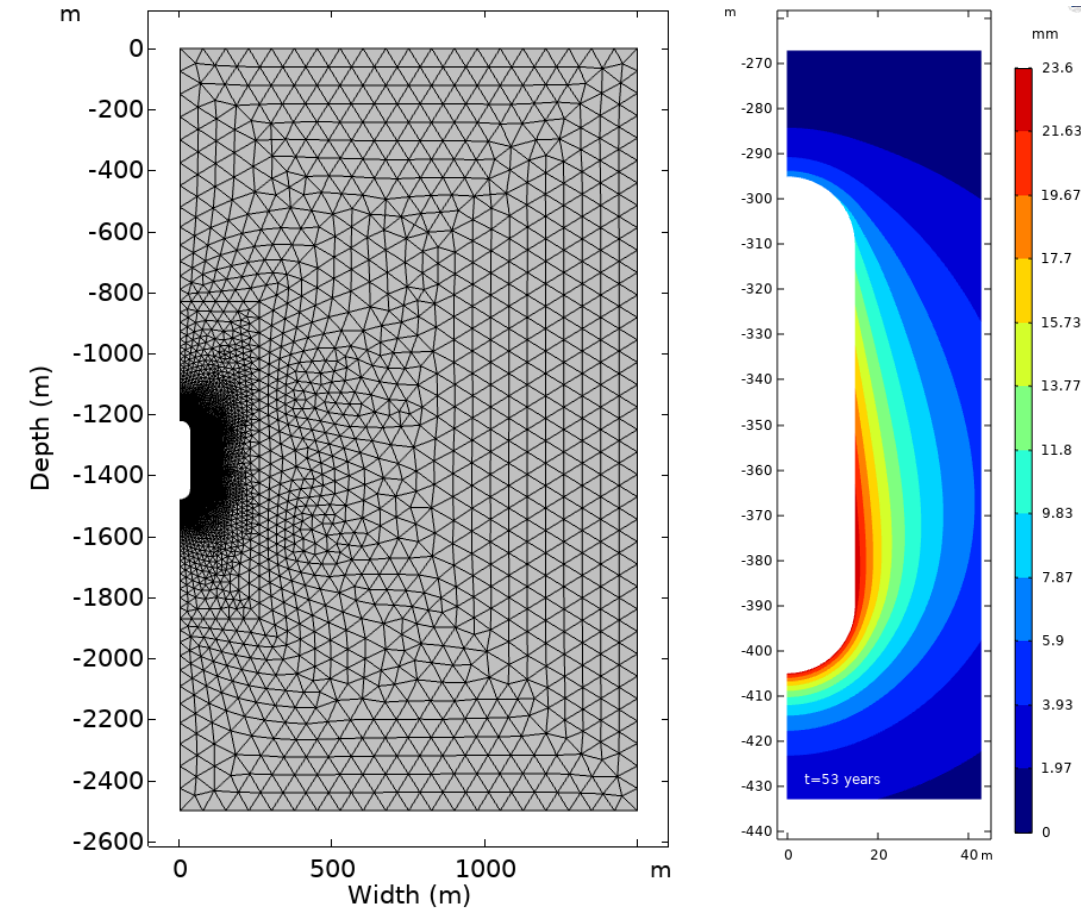
EXPERIMENTAL HIGHLIGHTS

- Permeability of the studied rock salt varies over more than 4 orders of magnitude
- Slippage (Klinkenberg) effect only observed for the less permeable and damaged samples
- Deviatoric loading under low confining pressure (1 MPa)
 - Small increase in gas permeability from the dilatancy threshold due to microcracking
- Deviatoric loading under high confining pressure (5 MPa)
 - No increase in permeability because the material becomes fully plastic (no more microcracking)
- Hydrostatic loading
 - Gas permeability decreases because of irreversible cracks closure due to self-healing process
- Permeability increases slightly during dynamic/cyclic mechanical and thermal fatigue tests (corresponding to high-frequency cycling in salt caverns) due to microcracking
- Permeability decreases during static fatigue (creep) thanks to the self-recovery process
- Results have to be confirmed on less initially damaged salt samples !!!

**The different mechanisms (viscoplasticity with strain hardening, microcracking and cracks healing) involved in rock salt deformation act in a competitive way to annihilate any significant permeability evolution.
→ Strong confidence in the H₂ storage in salt caverns which remains by far the SAFEST SOLUTION**

Numerical modeling of hydrogen storage in salt cavern

- From available experimental results on salt rock, propose/develop a **rheological model that reproduces the main features of short and long term behavior**
- **Validation based on experimental tests and application(s) to salt cavern**
- Analysis of the impact of the operational phase:
 - ✓ Mechanical behaviour of the salt cavern
 - ✓ H₂ leakage through the cavern wall (extent of the dissolved H₂ plume)



- **HydroMechanical model**: viscoplastic and damage model

Coarita-Tintaya E.D., Golfier F., Grgic D., Souley M. and Cheng L. 2023, Hydromechanical modelling of salt caverns subjected to cyclic hydrogen injection and withdrawal, Computers and Geotechnics, 162: 105690, [10.1016/j.compgeo.2023.105690](https://doi.org/10.1016/j.compgeo.2023.105690)

- **ThermoHydroMechanical model**: Previous model + fatigue + healing + thermal coupling

Elastoplastic and damage model

Elastic tensor

(compressive stresses are negative and $\sigma_1 \leq \sigma_2 \leq \sigma_3$)

$$C_{ijkl} = (1 - d) C_{ijkl}^0$$

$$d = d_i + d_f + d_t - h \quad ; \quad d \in [0; 1]$$

d_i : instantaneous damage
 d_f : fatigue damage
 d_t : tertiary damage
 h : healing

} **d: total damage**

Mohr-Coulomb yield surface

$$f_p = q + p M_p - N_p$$

$$M_p = \eta_p(\gamma^p) \frac{\sin \phi}{\left(\frac{\cos \theta}{\sqrt{3}} - \frac{1}{3} \sin \theta \sin \phi \right)} \quad ; \quad N_p = \eta_p(\gamma^p) \frac{(1 - d_i - d_f + h) c \cos \phi}{\left(\frac{\cos \theta}{\sqrt{3}} - \frac{1}{3} \sin \theta \sin \phi \right)}$$

Hardening rule

$$\eta_p(\gamma^p) = 1 - (1 - \eta_p^0) \exp \{ -\alpha_p^0 \gamma^p \}$$

Elastoplastic and damage model

Plastic potential

(Chiarelli *et al* 2003,
Souley *et al.* 2017)

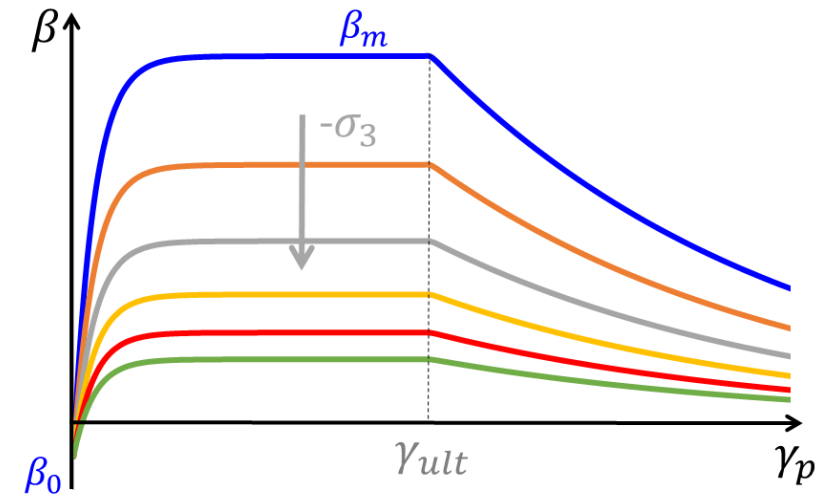
$$g = q + \beta(\gamma^p) p$$

$$\beta(\gamma^p) = \begin{cases} \beta_m - (\beta_m - \beta_0) \exp(-b_\beta \gamma^p) & ; \gamma^p < \gamma_{ult}^p \\ \beta_{ult} \exp\left(1 - \frac{\gamma^p}{\gamma_{ult}^p}\right) & ; \gamma^p \geq \gamma_{ult}^p \end{cases}$$

$$\beta_m = \begin{cases} \beta_m^0 \exp(a_p \sigma_1) & ; \sigma_1 < 0 \\ \beta_m^0 & ; \sigma_1 \geq 0 \end{cases}$$

$\beta(\gamma^p)$: dilatancy coefficient

β_0 ($\beta_0 < 0$) is the initial volumetric contraction, β_m ($\beta_m > 0$) is the volumetric dilatancy at large deformation



We have volumetric contraction when $\beta < 0$, whereas the plastic strains evolve towards volumetric dilatancy if $\beta > 0$.

Elastoplastic and damage model

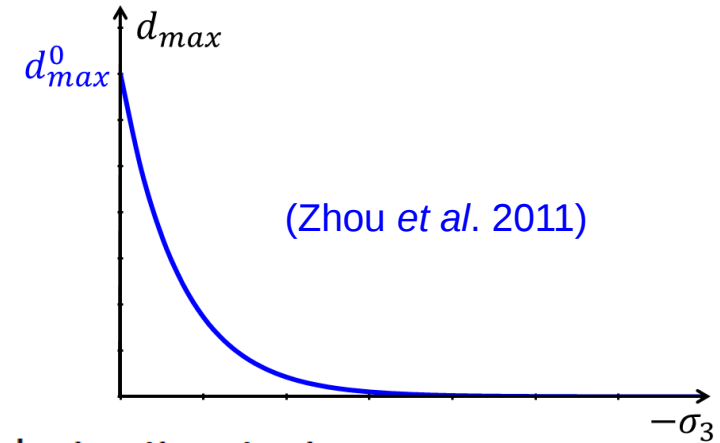
Instantaneous damage (d_i) (Mazars 1984)

$$f_d = d_i^{max} \{1 - \exp[-b_d (Y_d - Y_0)]\} - d_i \leq 0 \quad ; \quad d_i \in [0; d_i^{max}]$$

$$d_i^{max} = \begin{cases} d_{i0}^{max} \exp(a_d \sigma_1) & ; \quad \sigma_1 < 0 \\ d_{i0}^{max} & ; \quad \sigma_1 \geq 0 \end{cases}$$

$$Y_d = \sqrt{\langle \varepsilon \rangle : \langle \varepsilon \rangle} ; \varepsilon \text{ is the principal strain value of the elastic and plastic strains}$$

a_d : parameter describing the effect of confining stress



Fatigue damage (d_f) (Zhang et al. 2023)

$$\dot{d}_f = \frac{b_f}{T_f} (d_f^{max} - d_f) \quad ; \quad d_f \in [0; d_f^{max}]$$

$$d_f^{max} = d_{f0}^{max} \eta_p \left(1 + b_{f2} \frac{\langle Y_d - Y_d^c \rangle}{Y_d^c} \right)$$

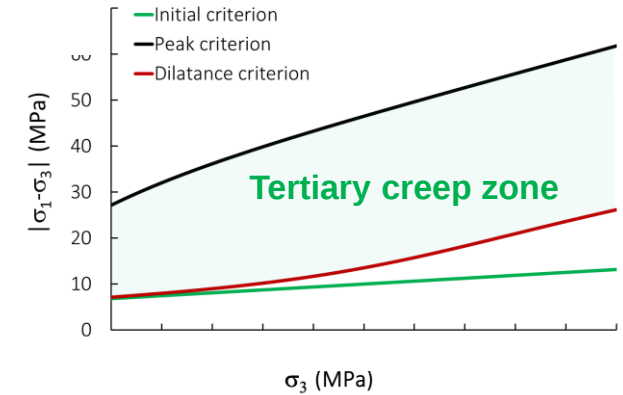
Assumption: The value of instantaneous damage is constant during the fatigue test

Elastoplastic and damage model

Creep strain rate Transient creep Steady-state creep
(Kelvin-Voigt model) (Norton law)

$$\dot{\epsilon}_{ij}^c = \left[\frac{1}{\bar{\eta}_k^*} (q^* - \bar{G}_k^* \epsilon_{tr}) + A_N \exp\left(-\frac{B_N}{T}\right) \left(\frac{q^*}{\sigma_{ref}}\right)^{n_N} \right] \frac{3}{2} \frac{s_{ij}}{q}$$

$$\bar{G}_k^* = G_k \exp(k_1 q^*) ; \bar{\eta}_k^* = \eta_k \exp(k_2 q^*)$$



Tertiary damage (d_t)

(Hou *et al.* 2003)

$$\dot{d}_t = A_T \frac{A_T}{(1-d)^{n_T}} \left\langle \frac{f_{ds}}{\sigma_{ref}} \right\rangle^{n_T}$$

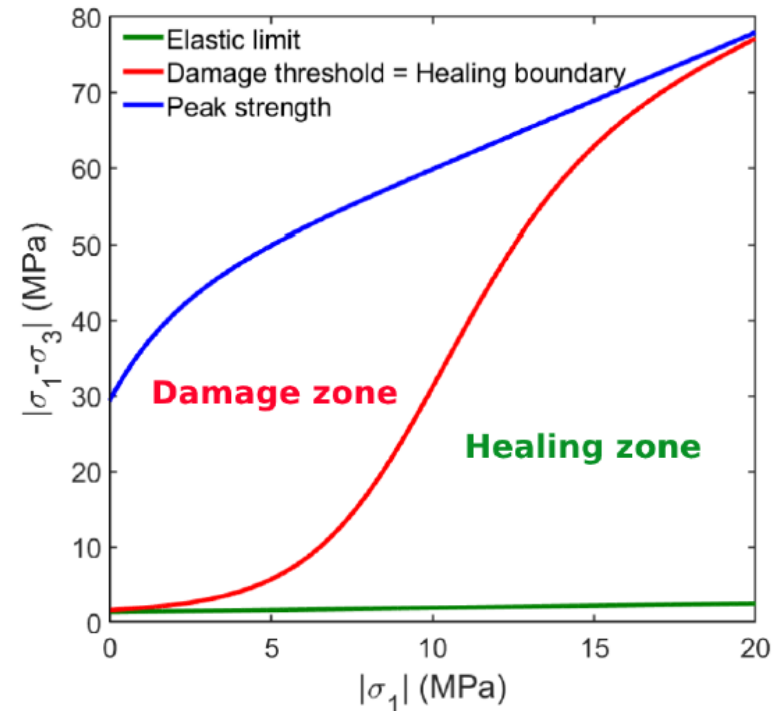
$$q^* = \frac{q}{1-d} \quad \begin{array}{l} q^*: \text{Mises stress undamaged} \\ q : \text{Mises stress damaged} \end{array}$$

→ Damage is activated and increases if the damage limit (dilatance criterion) is exceeded

Healing (h)

$$\dot{h} = \frac{d}{h_1} \left\langle \frac{-f_{ds}}{\sigma_{ref}} \right\rangle$$

Assumption: The healing boundary and the damage threshold are identical



Mathematical formulation

Thermo-hydro-mechanical coupling

$$\nabla \cdot [C : \varepsilon^e - b(p - p_{ref}) \mathbf{I} - \alpha C : (T - T_{ref}) \mathbf{I}] + \rho_m \vec{g} = 0$$

$$\rho_f b \frac{\partial \varepsilon_v}{\partial t} + \frac{\rho_f}{M} \frac{\partial p}{\partial t} + \nabla \cdot (\rho_f \vec{q}) - 3\rho_f \alpha_m \frac{\partial T}{\partial t} = 0$$

$$\rho C_p \frac{\partial T}{\partial t} + \cancel{\rho_f C_{p,f} \vec{q} \cdot \nabla T} - \nabla \cdot (\lambda \nabla T) = 0$$

Where:

$$n \frac{\partial (c_{H_2})}{\partial t} + \nabla \cdot (q_i c_{H_2}) = \nabla \cdot (n \mathbf{D}^* \cdot \nabla c_{H_2})$$

$$C = (1 - d) C^0$$

$$\rho_m = (1 - n) \rho_s + n \rho_f$$

$$\vec{q} = -\frac{k}{\mu_f(T)} (\nabla p + \rho_f \vec{g})$$

$$\frac{1}{M} = \frac{n}{K_f} + \frac{(1 - b)(b - n)}{K_d}$$

$$\alpha_m = (b - n) \alpha + n \alpha_f(T)$$

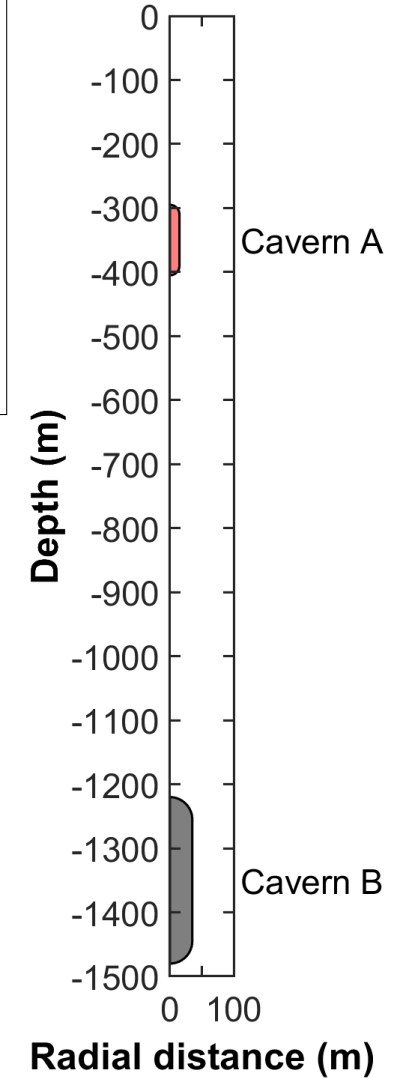
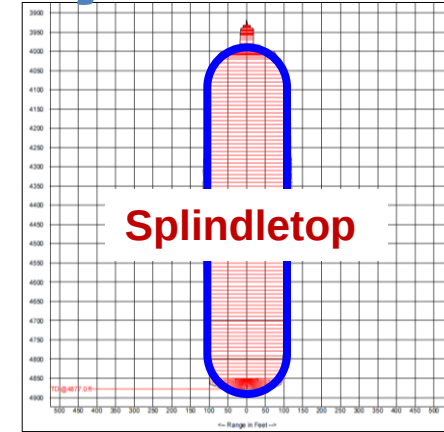
$$k = k_0 10^{A_k d} \quad (\text{Gawin et al. 2002})$$

- Fully water-saturated material
- Two-phase flow not considered
- Dissolved H_2 transport by advection and diffusion
- The value of the Biot coefficient is 0.3 if material is damaged

Storage in salt cavern : case study

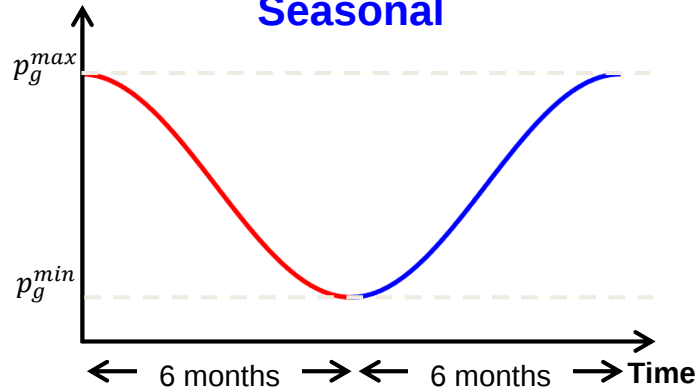
boundary and initial conditions

			Cavern A	Cavern B
Model geometry	Depth	m	-350	-1350
	Volume	m ³	70 686	910 800
H ₂ Conditions	P _{max} - H ₂	MPa	5	20
	P _{min} - H ₂	MPa	2	6



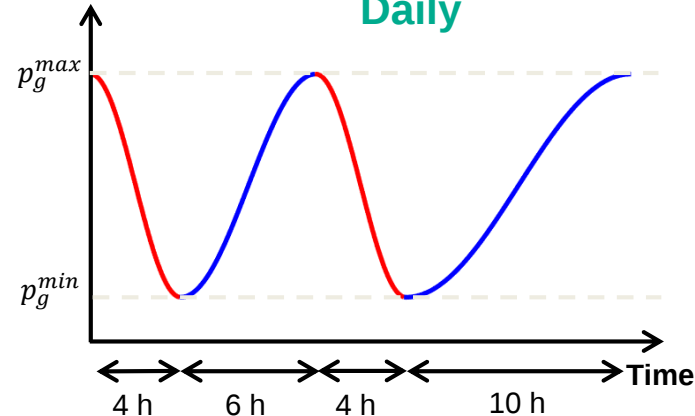
Cycling scenarios:

Seasonal



→ For 50 years

Daily



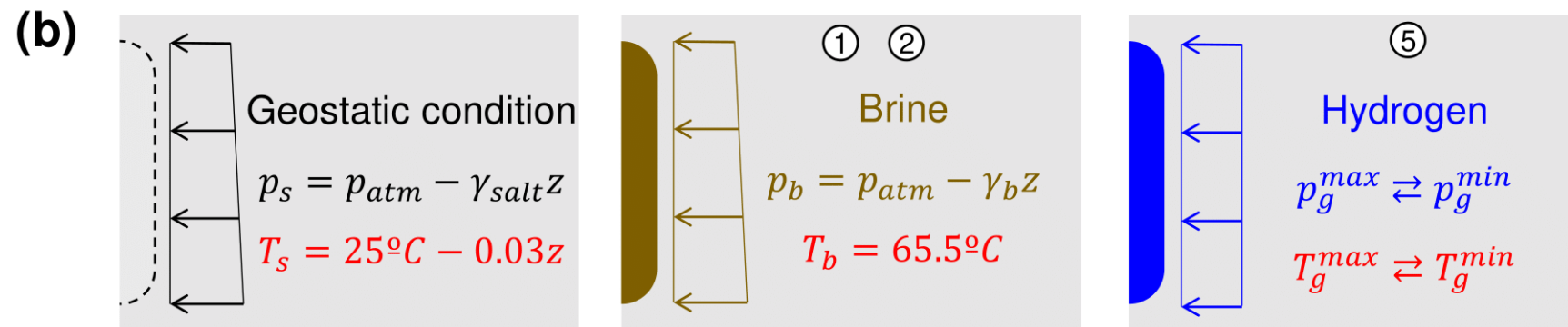
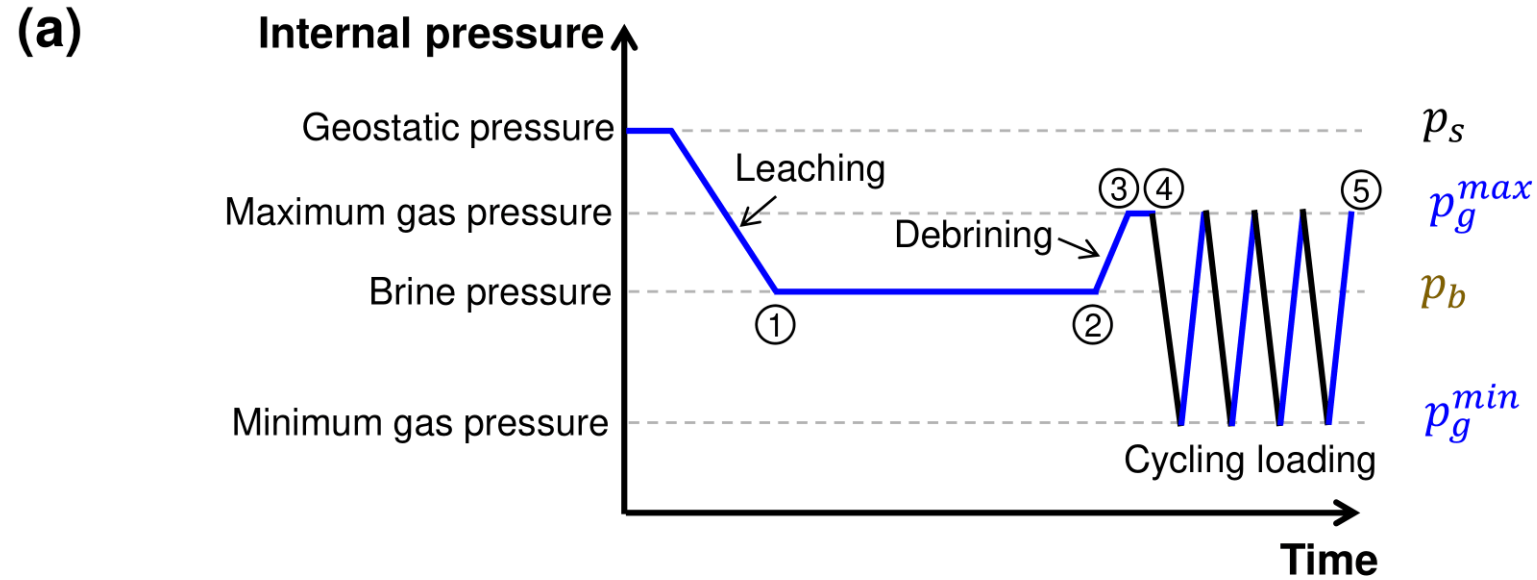
→ For 100 days

Analysis cases

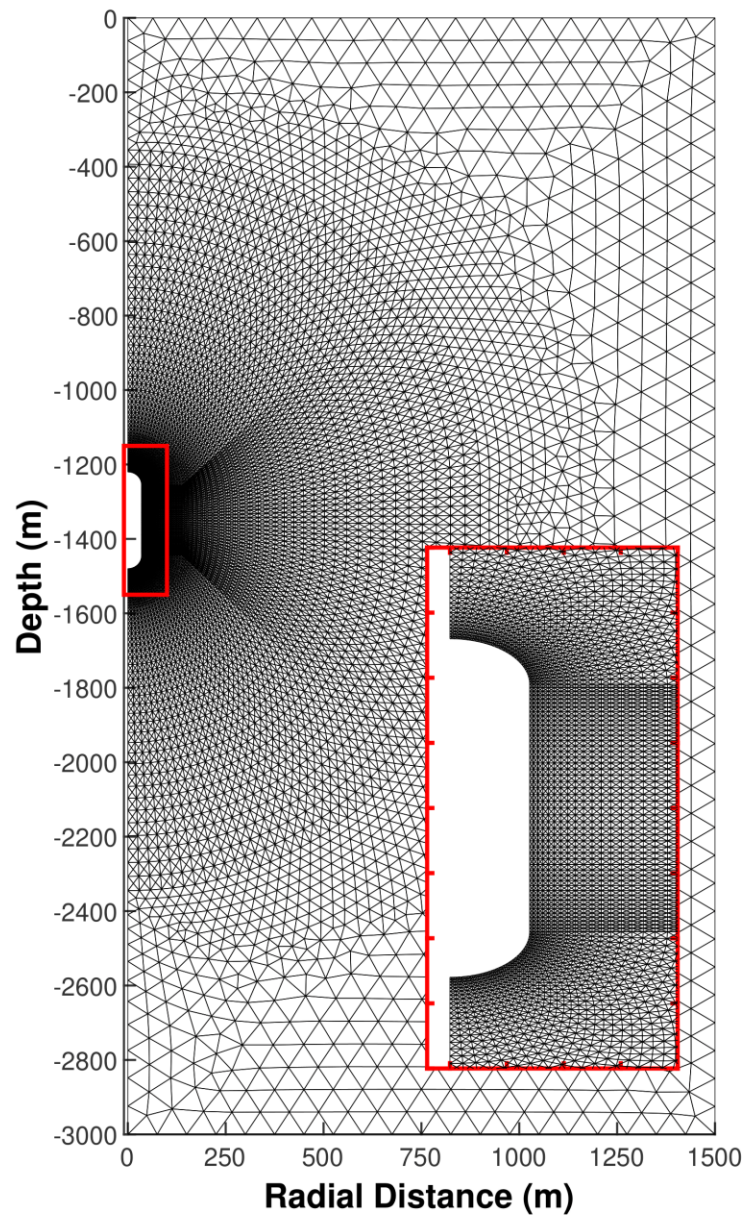
- Elastic (→ To verify the HM model: realised)
- Cavern A (seasonal & daily)
- Cavern B (seasonal & daily)

Storage in salt cavern : case study

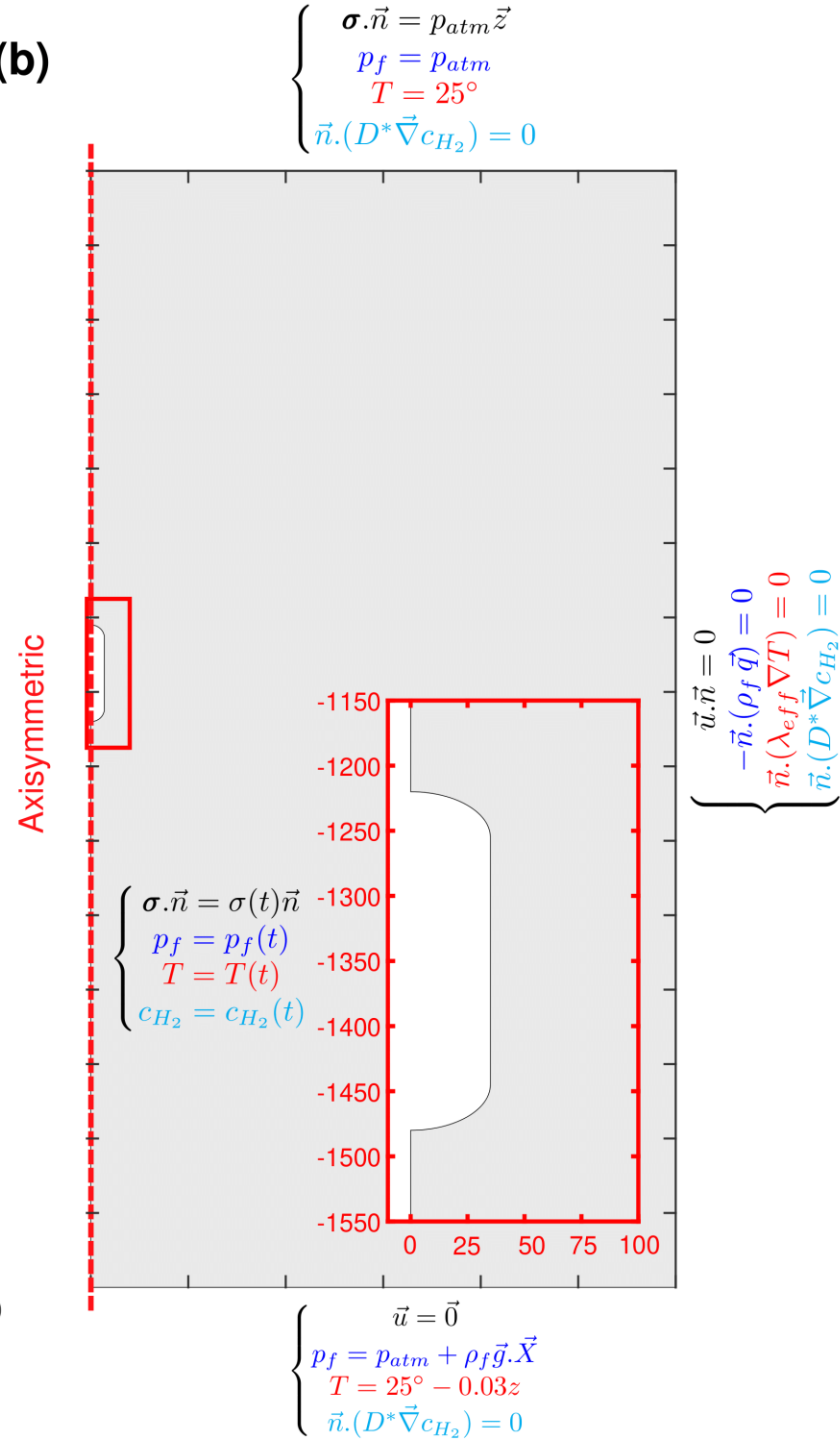
Cavern geometry, boundary and initial conditions and properties



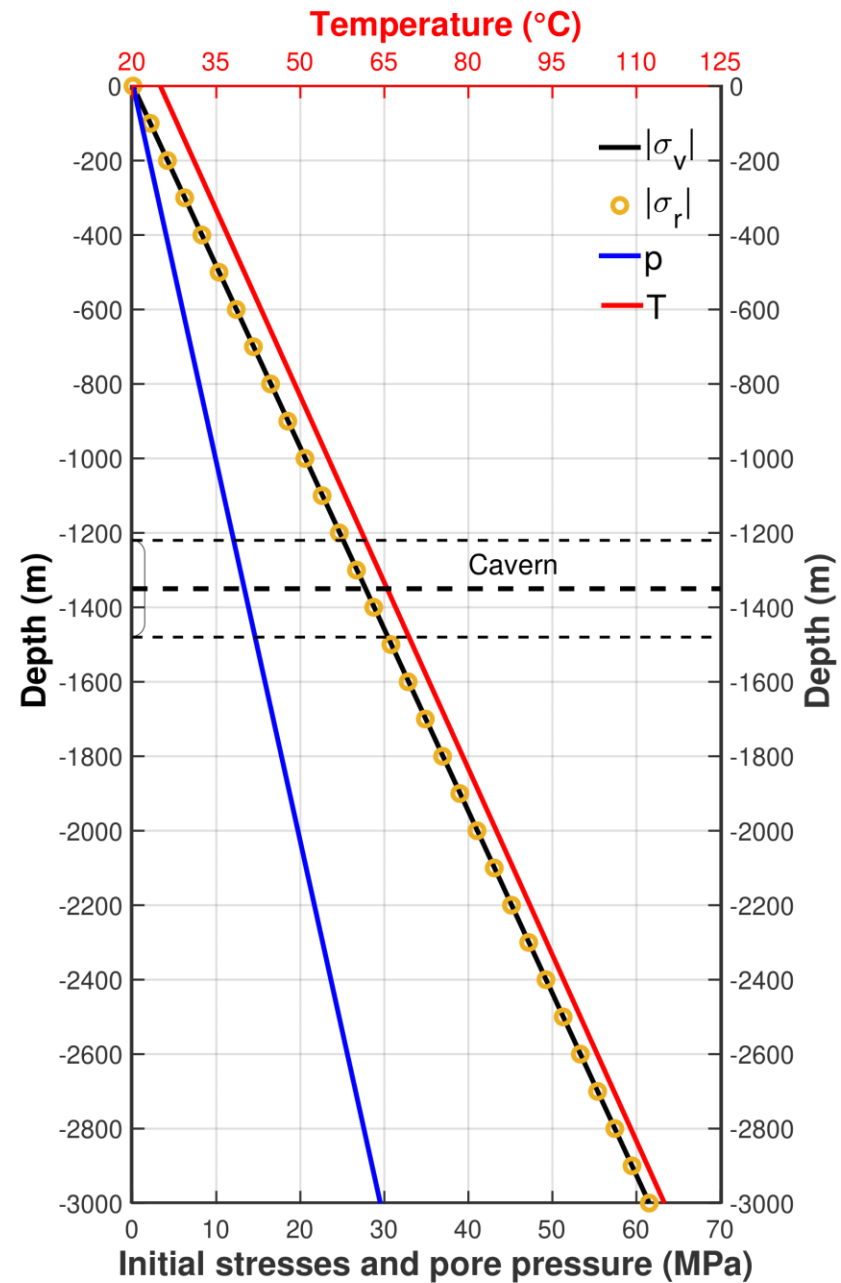
(a)



(b)



(c)



Storage in salt cavern : case study

Cavern geometry, boundary and initial conditions and properties

Short-term numerical values

Parameter	Value	Unit	Parameter	Value	Unit	Parameter	Value	Unit	Parameter	Value	Unit
E	35.0	GPa	α_p^0	70		a_p	0.5	1/MPa	d_{f0}^{max}	0.031	
ν	0.3		β_0	-0.02		b_d	70		b_f	0.036	s
c	12.5	MPa	β_m^0	1.5		d_{i0}^{max}	0.3		b_{f2}	1	
ϕ	28.3	°	b_β	80		a_d	0.5	1/MPa	Y_d^c	0.0278	
η_p^o	0.05		γ_{ult}^p	0.09		R_t	3	MPa	T_f	1	y

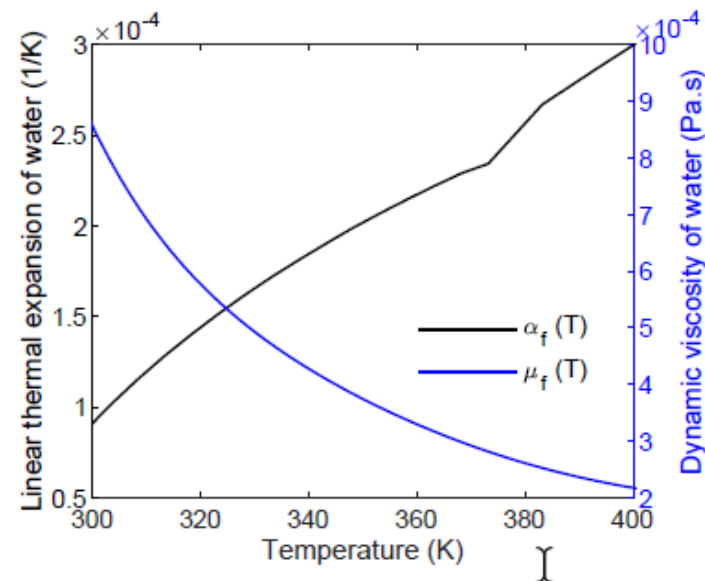
Long-term numerical values

Parameter	Value	Unit	Parameter	Value	Unit	Parameter	Value	Unit
G_k	12600	MPa	η_k	60000	MPa.d	n_T	3.5	
k_1	-0.18	1/MPa	k_2	-0.15	1/MPa	A_d	4.0E-10	1/d
l_1	0	1/K	A_N	0.005	1/d	β_{d0}	1	
G_{kE}	7560	MPa	B_N	4700	K	a_t	0.5	1/MPa
k_{1E}	-0.18	1/MPa	n_N	4		h_1	6000	d
l_{1E}	0	1/K	A_T	4.0E-8	1/d	β_h	1	

Storage in salt cavern : case study

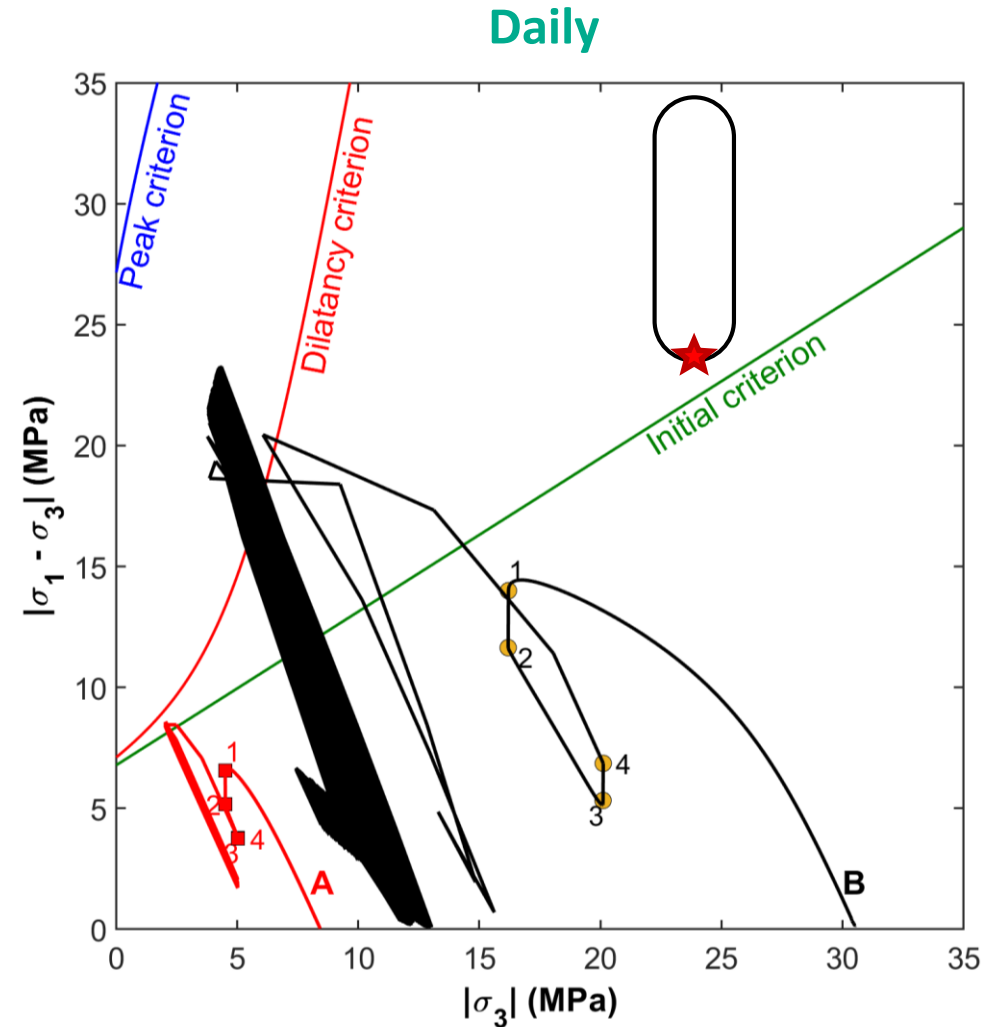
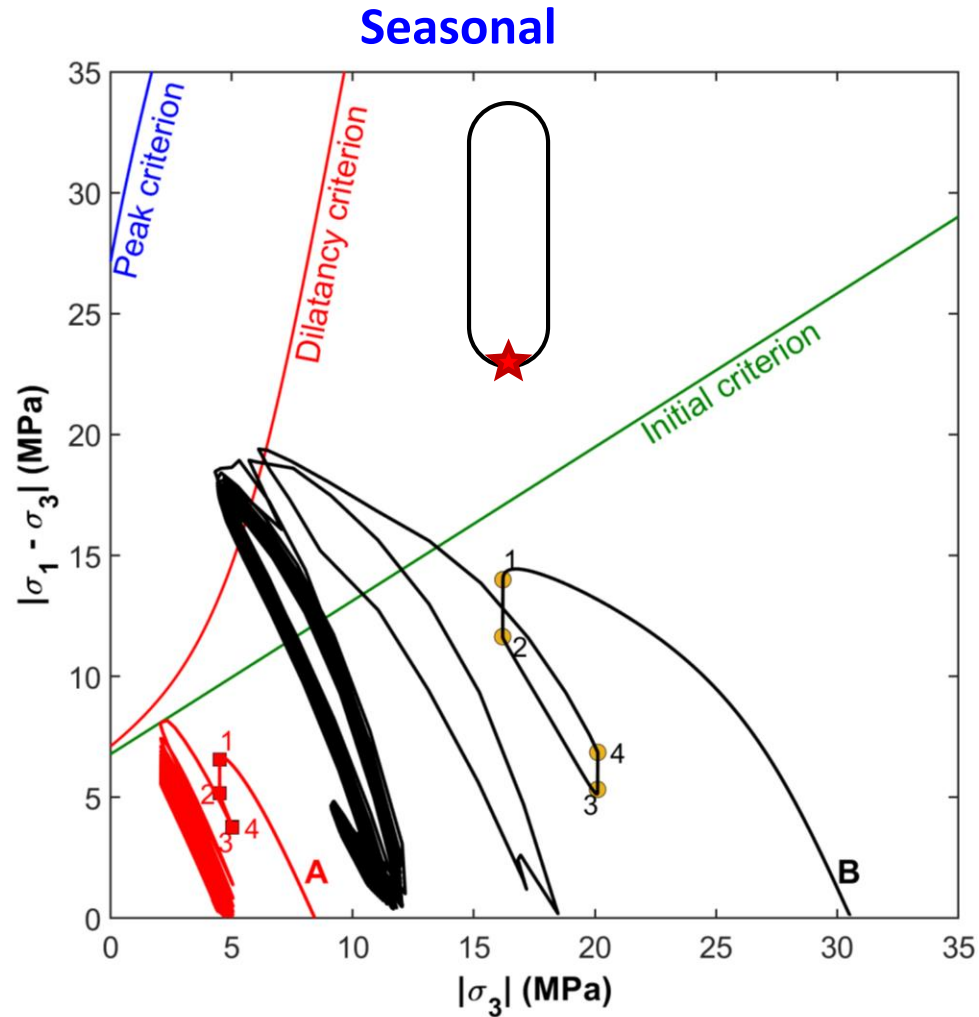
Cavern geometry, boundary and initial conditions and properties

Hydraulic			Thermal			Transport		
Parameter	Value	Unit	Parameter	Value	Unit	Parameter	Value	Unit
k_0	1.0E-20	m ² /s	λ	6.8	W/(m.K)	D^*	6.0E-9	m ² /s
A_k	90		α	4.0E-5	1/K			
μ_f	0.001	Pa.s	C_p	850	J/(kg.K)			
n	0.012		α_f	$f(T)$	1/K			
b	0.012							
K_f	2.2E+9	Pa						
ρ_f	1000	kg/m ³						



Hydro-mechanical results

Stress path at the bottom of the cavern (Deviatoric stress projected to the plane $\theta=\pi/6$)

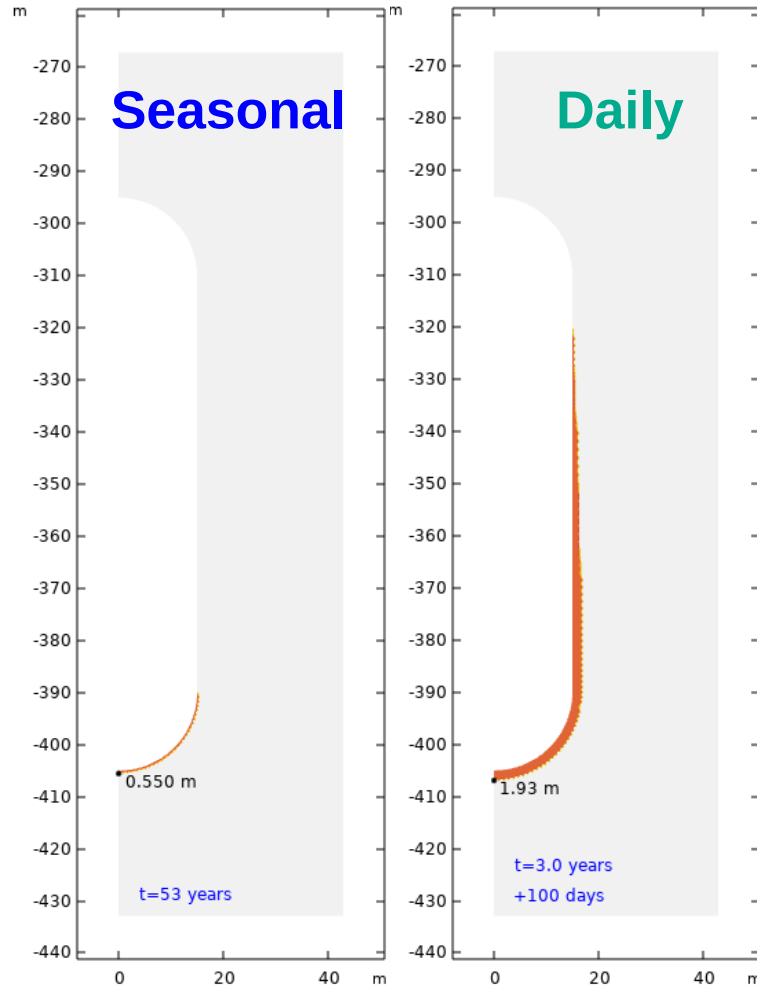


- Relaxation of deviatoric stresses over time
- **Seasonal and cavern A:** in the first extraction it exceeds the elastic limit. For the following cycles this limit is not exceeded due to creep. In **cavern B** the dilatancy criterion is reached

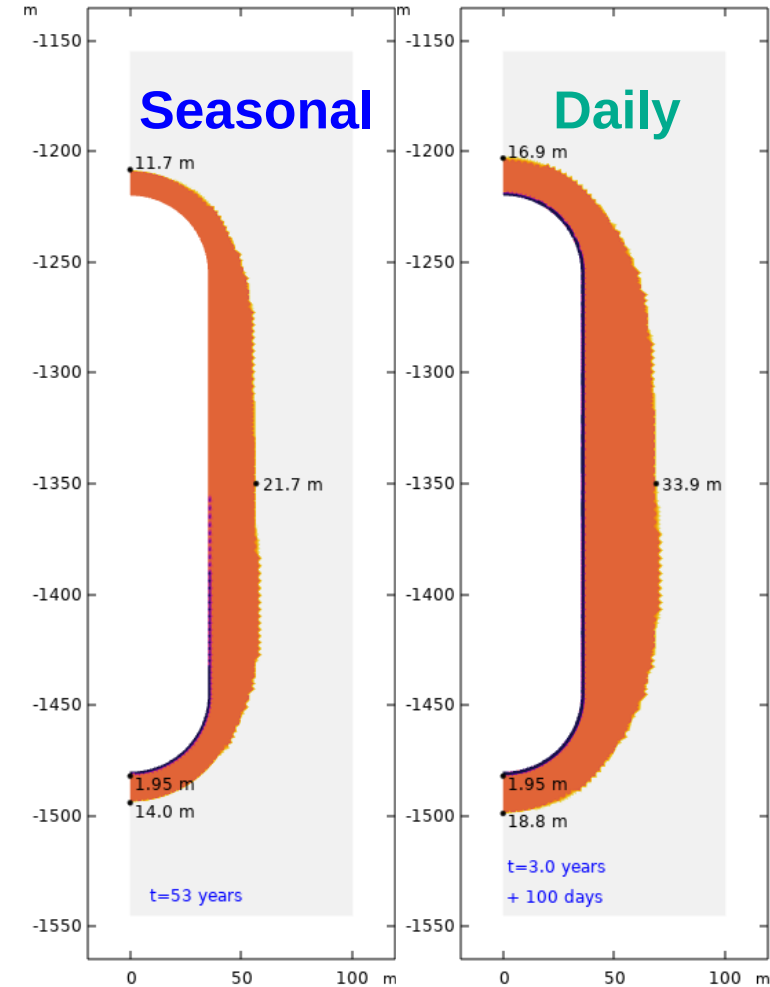
Hydro-mechanical results

Plastic and damage zone

Cavern A



Cavern B



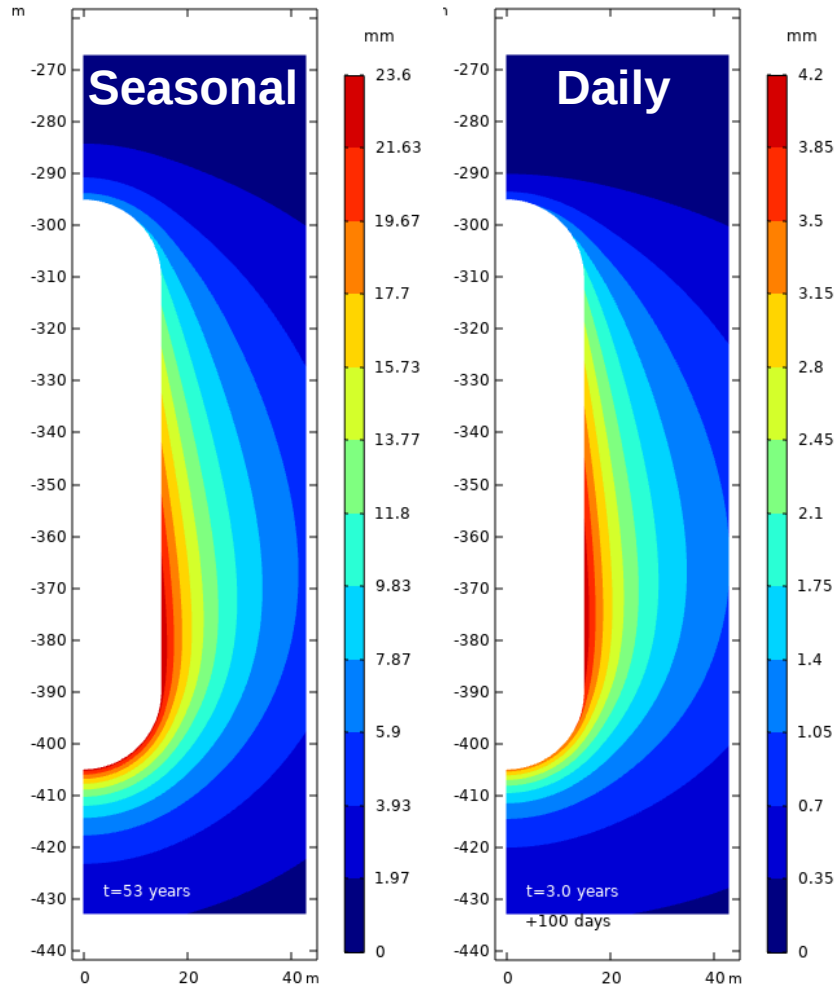
Around **cavern A** there is a small extension of the plastic zone, in the lower and lateral part of the cavern.

Around **cavern B**, the damage zone is initially located at the bottom. The plastic zone is all around the cavern and is more extensive for daily cycling

Hydro-mechanical results

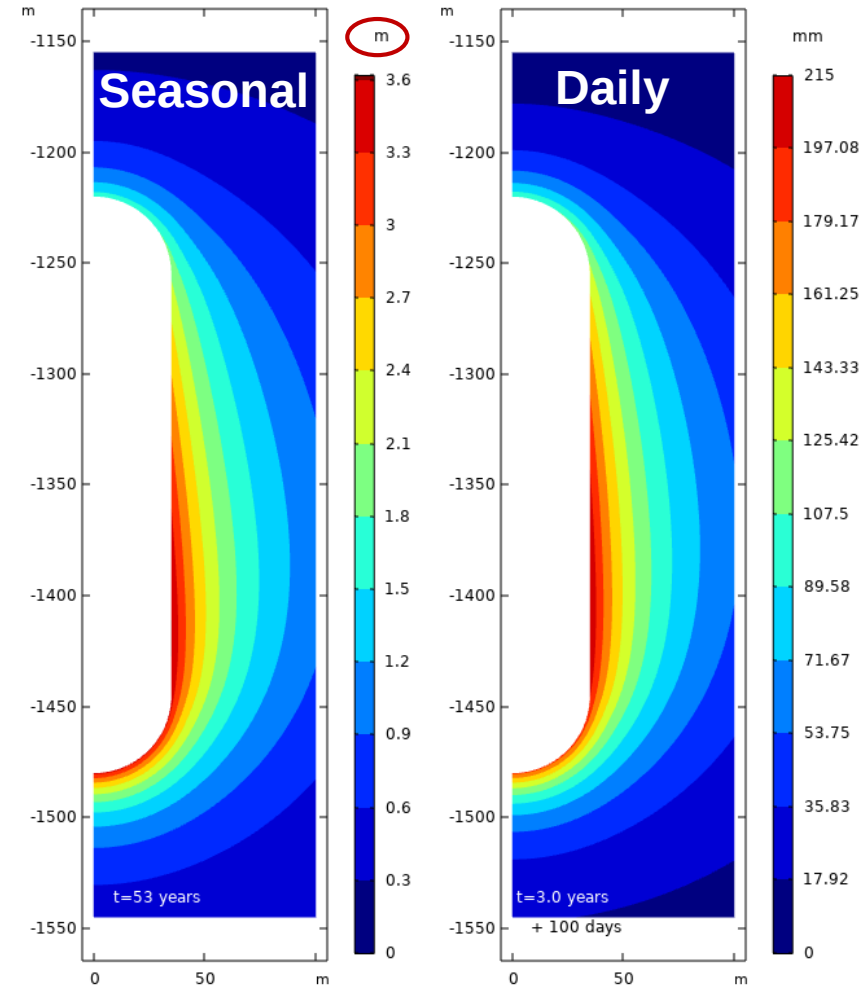
Cavern wall displacements

Cavern A



Around **cavern A**, the displacements are small and in the order of 26 mm

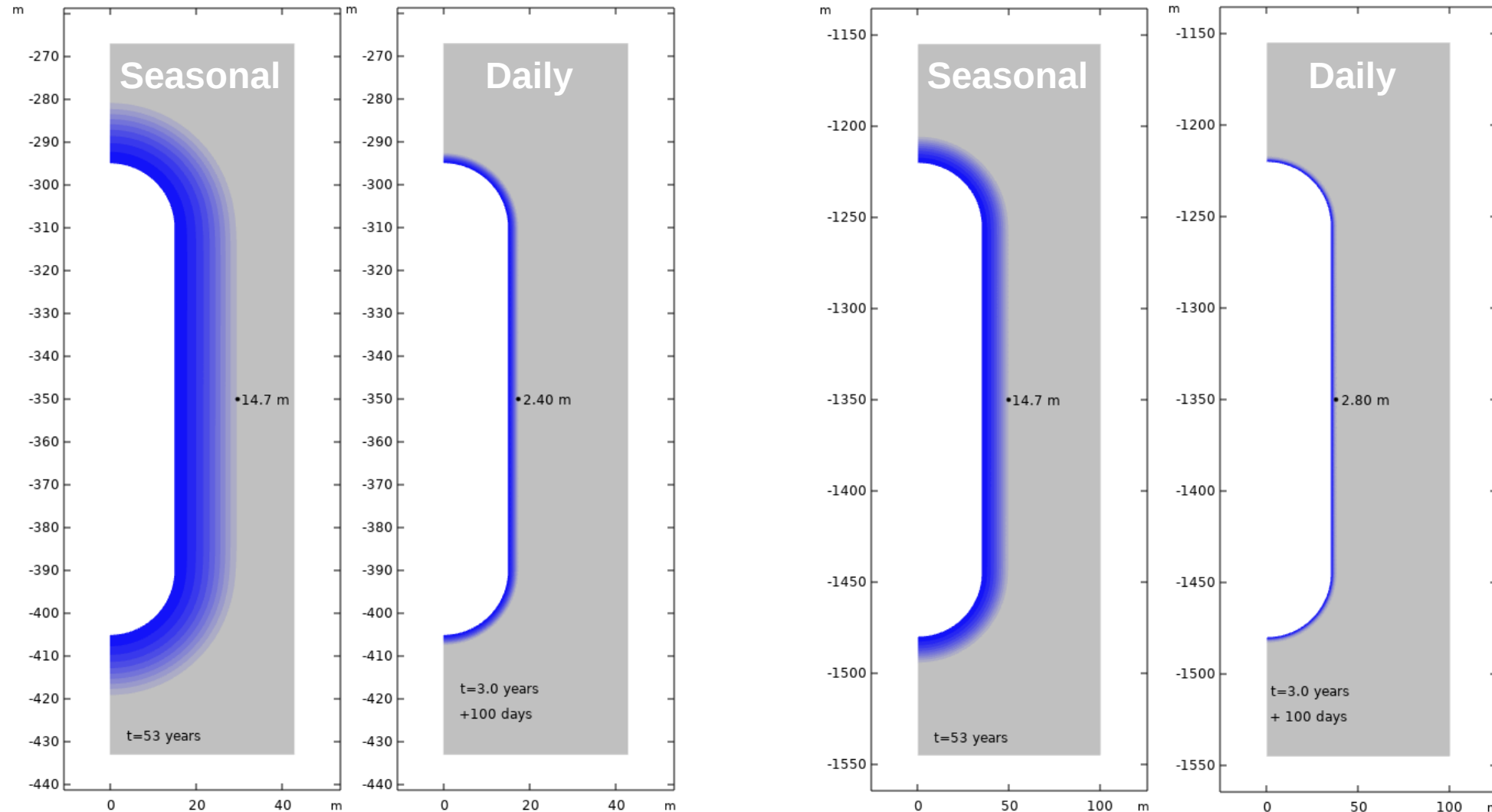
Cavern B



Around **cavern B**, the displacements are important for seasonal cyclage. Maximum displacement of 3.6 m on the lateral side

Hydro-mechanical results

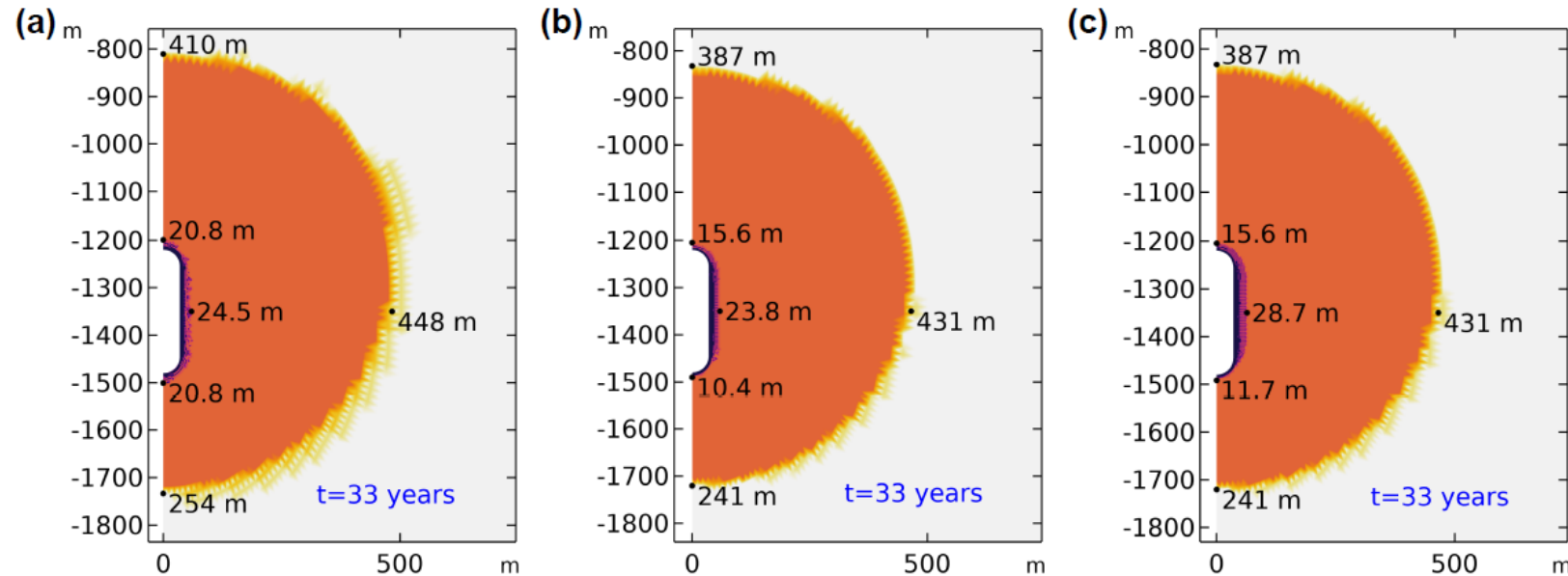
Hydrogen leakage



For both caverns, a similar hydrogen extension is estimated around the cavern. Almost 2.5 m for a daily cycling and almost 15 m for a seasonal cycling. **Gas transport mainly by diffusion.**

Thermo-hydro-mechanical results

Plastic and damage zones around cavern

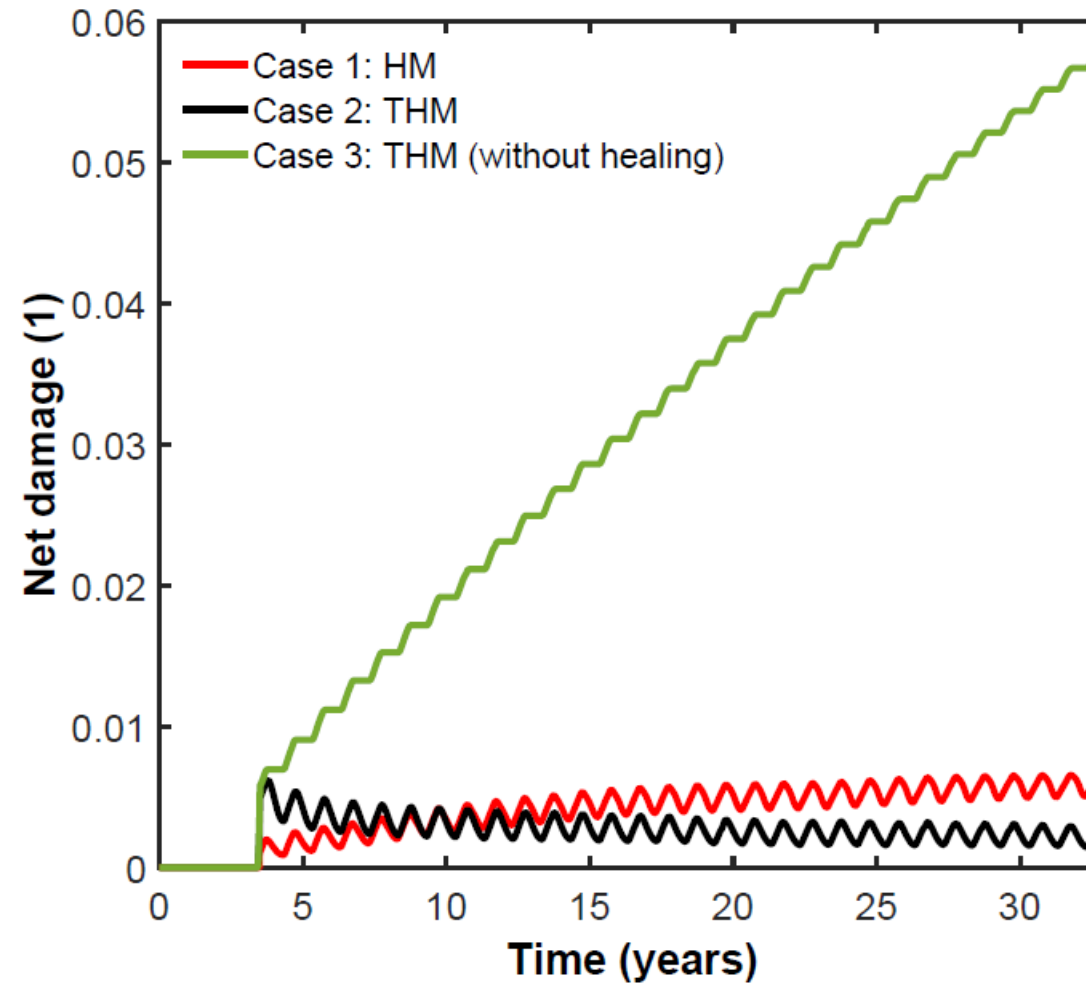


Zones at 33 years for (a) **Case 1**: HM (b) **Case 2**: THM (c) **Case 3**: THM (without healing)

- Before the operation phase, only plastic zones are calculated in the analysis cases
- At 3.5 years (first minimum gas pressure), the damage zone appears, and the plastic zone increases
- In the following years, damage zones increased slightly in extension, while the plastic zone increased more

Thermo-hydro-mechanical results

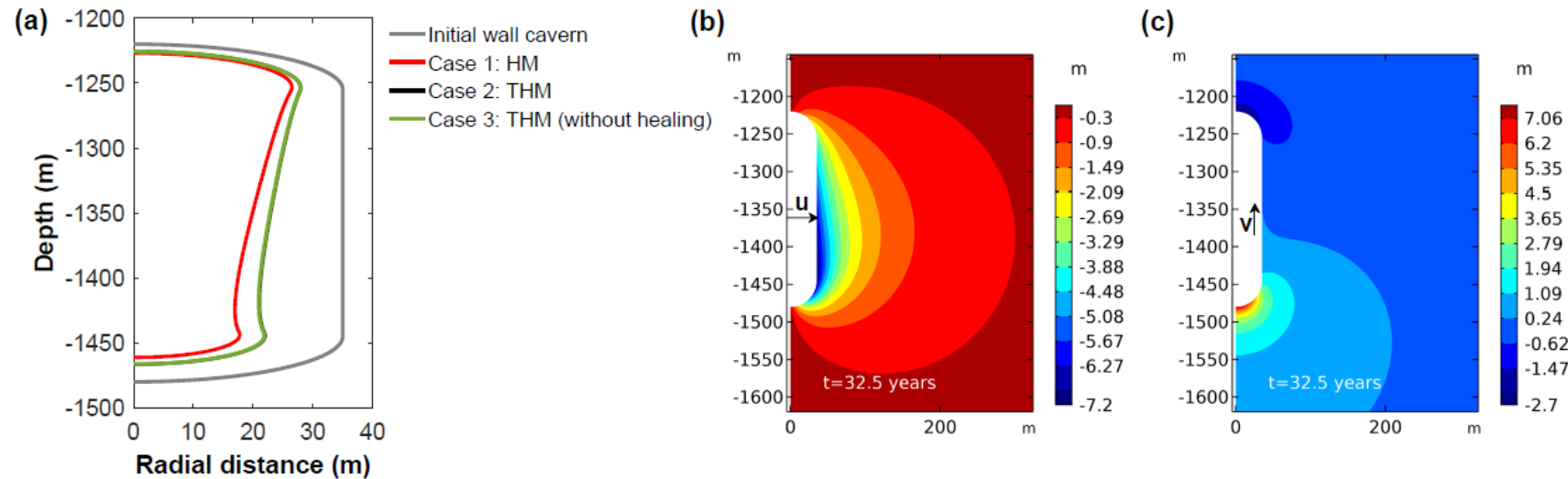
Net damage at the bottom of the cavern



Note: Net damage is defined as $d = d_i + d_f + d_t - h$

Thermo-hydro-mechanical results

Cavern wall displacements

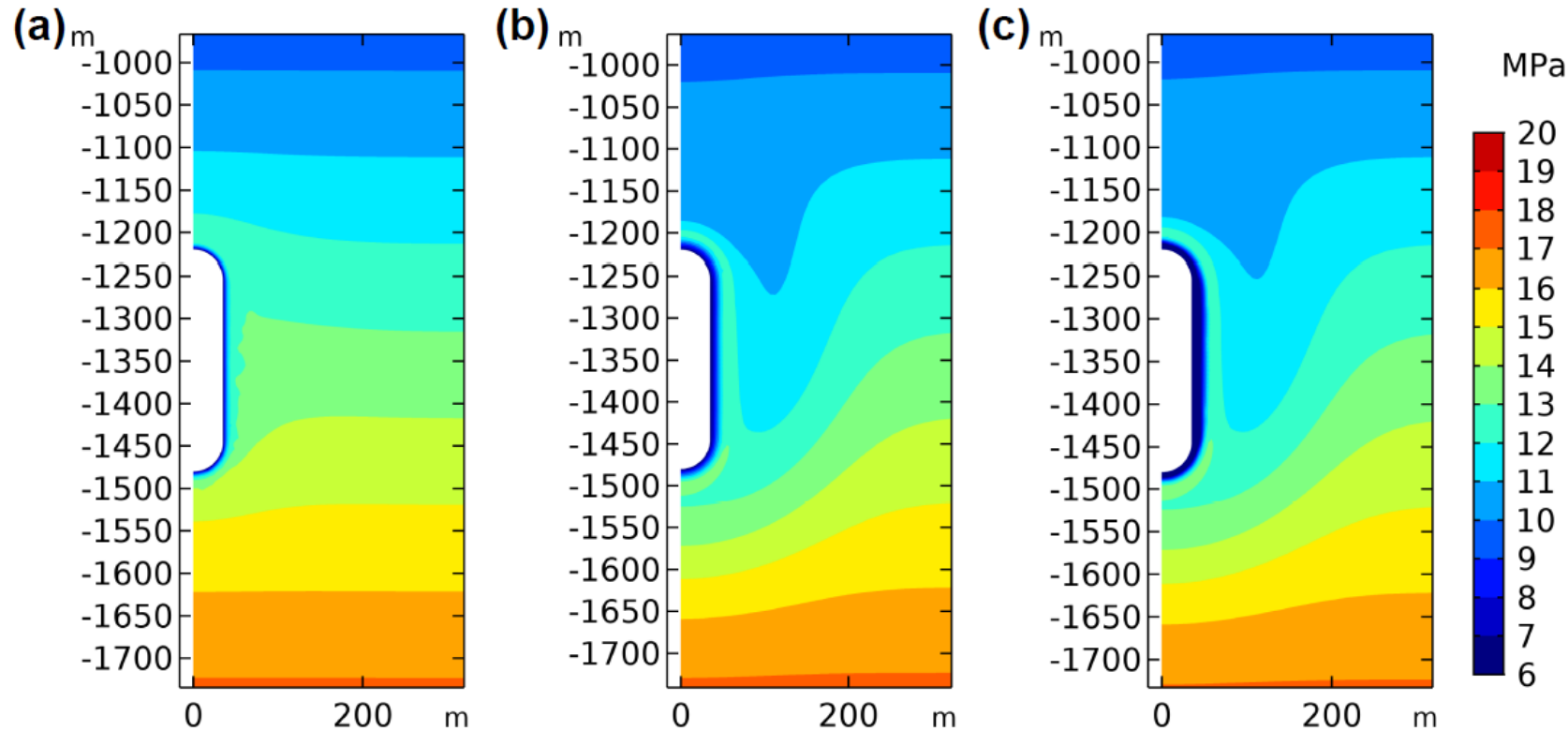


(a) Cavern wall displacements (scaled by a factor of 2.5), (b) horizontal and (c) vertical displacements for **Case 1: HM**

- Case 1 has more displacements than Cases 2 and 3 because it has a larger extent of the plastic zone around the cavern
- Cases 2 (THM) and 3 (THM without healing) exhibit similar displacements
- Important horizontal displacements are calculated on the bottom wall of the cavern, while significant vertical displacements are calculated on the floor of the cavern

Thermo-hydro-mechanical results

Pore pressure distribution

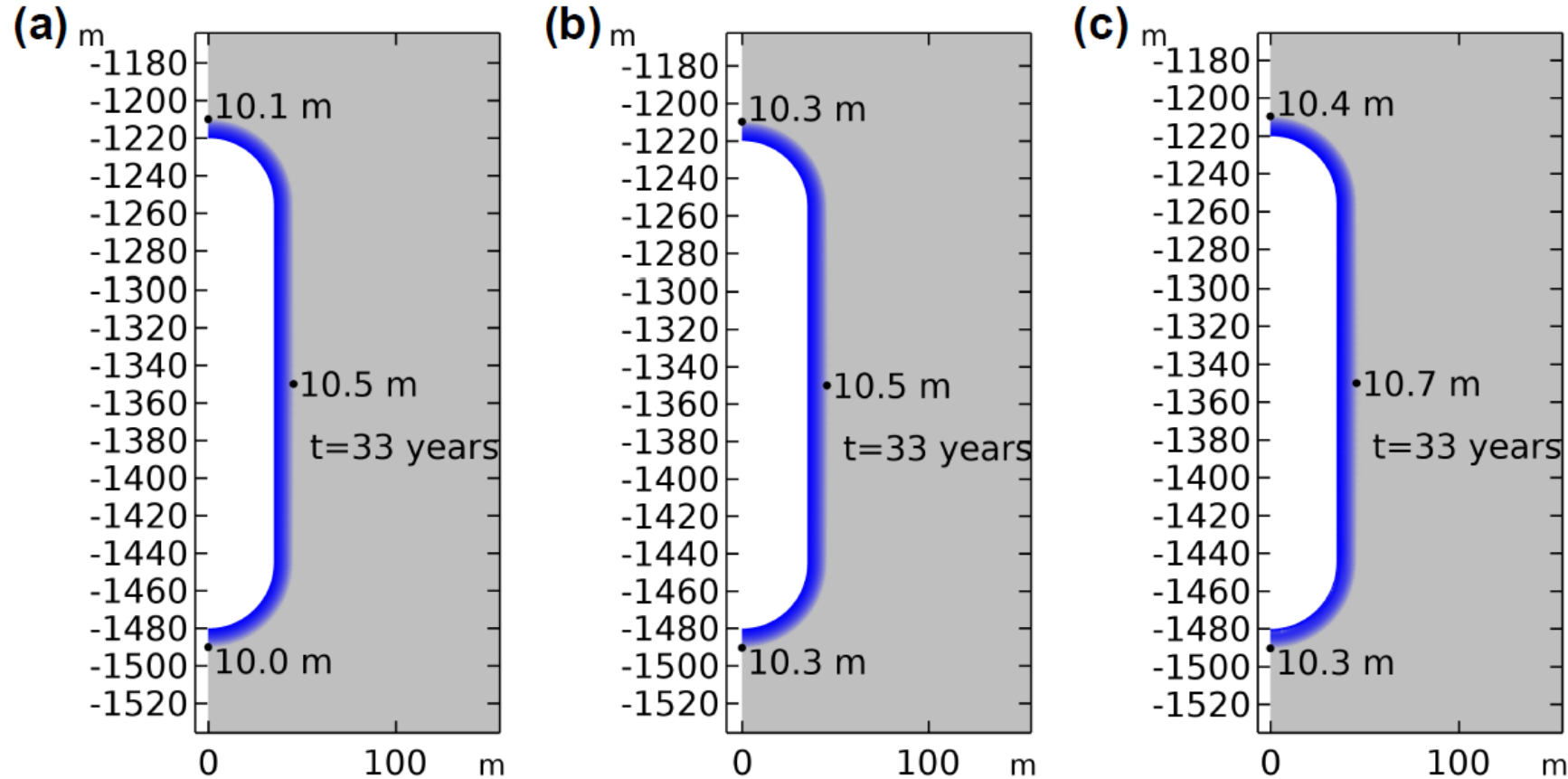


Pore pressure distribution around salt cavern at 32.5 years for:
(a) **Case 1**: HM (b) **Case 2**: THM (c) **Case 3**: THM (without healing)

- Pore pressure distributions for Cases 2 and 3 show more modification around the cavern compared to Case 1, due to the temperature effect.

Thermo-hydro-mechanical results

Hydrogen leakage



Hydrogen extension around salt cavern at 33 years for:

(a) **Case 1**: HM (b) **Case 2**: THM (c) **Case 3**: THM (without healing)

- Hydrogen extensions are similar for the three cases due to transport properties, modified only near the cavern wall

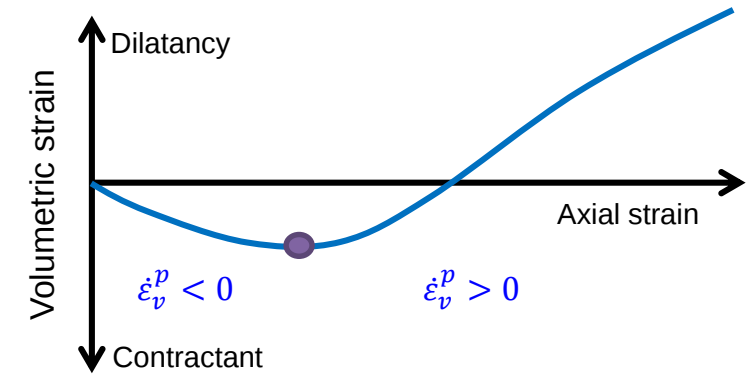
Conclusions

- A THM model describing the main key features of the rock salt behavior has been developed
 - ✓ **Short term model** takes into account elastoplastic and instantaneous damage behaviors
 - ✓ For the **long-term behavior**, the three creep phases generally observed on creep tests are considered
 - ✓ **Fatigue damage** and **healing** have been also implemented
 - ✓ **Thermal coupling** and **hydrogen transport** are considered
- **Numerical simulation results**
 - The **deep cavern** is more susceptible to mechanical stability problems.
 - The **daily scenario** is also **more detrimental** to the stability.
 - Even if damage occurs, the extent of damage zone is limited
 - the **amount of gas leakage is also limited** and both caverns lead to almost the same plume size because of diffusion which is preponderant.
 - Healing contributes to **annihilate any significant damage evolution**

Thank you for your attention !

Main assumptions of the model

- Rock salt is considered as an isotropic material
- Its elastic limit is very low -> The initial yield strength is assumed to be a fraction of the peak strength criterion
- It has a **hardening deformation mechanism** and shows a more **ductile response** than most other rocks
- **Ductile behaviour:** increasing confining stress
- Significant dilatancy at low confining stresses -> **To define the dilatancy criterion**
- we use the plastic potential: volumetric plastic strain (ε_v^p)
- if $\dot{\varepsilon}_v^p > 0$: we suppose that the material is in volumetric dilatancy

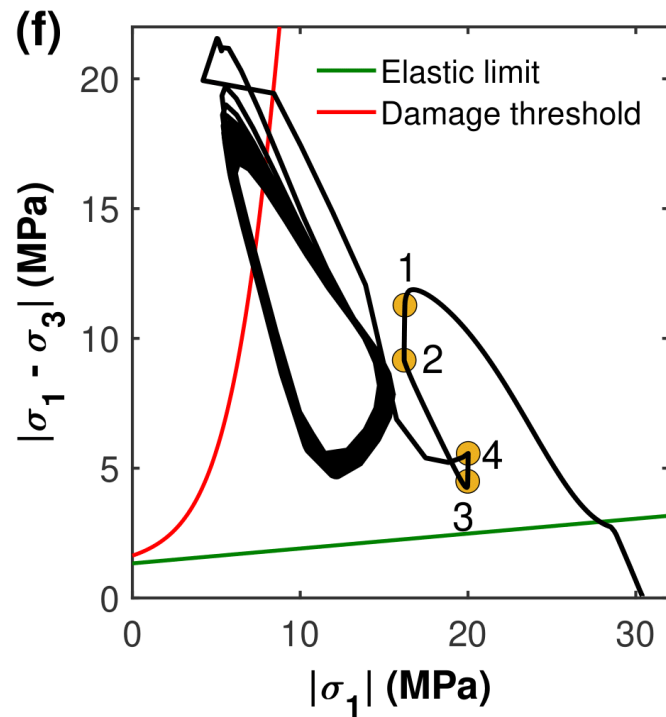
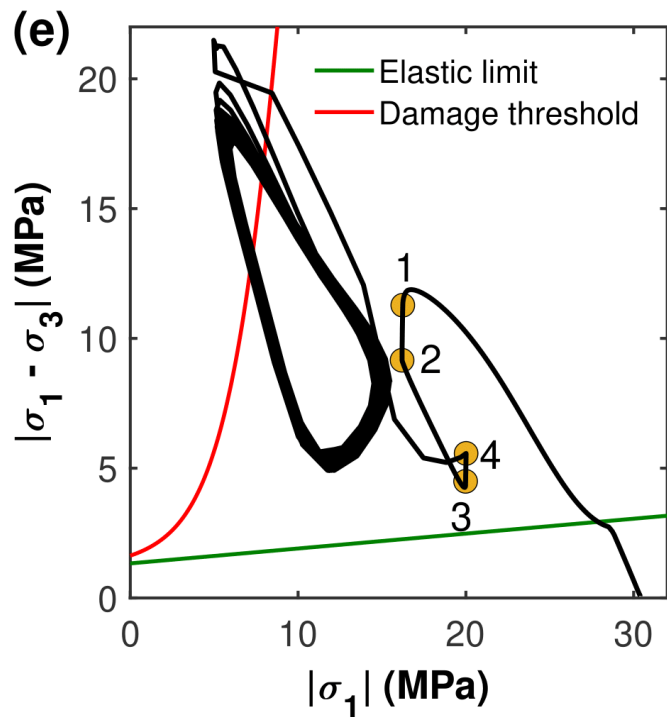
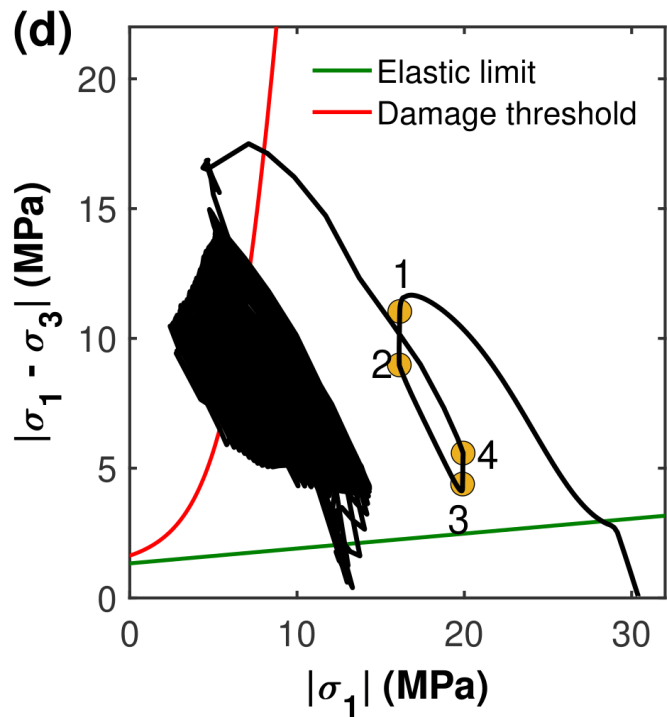
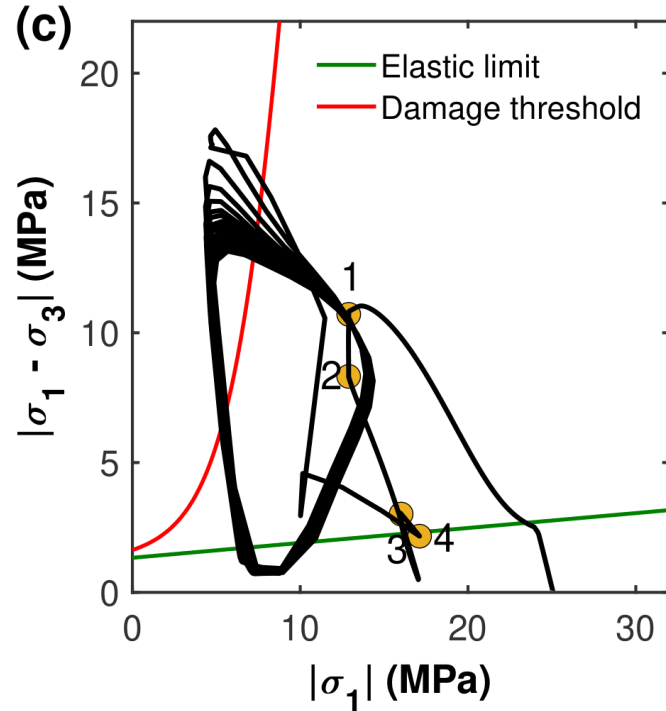
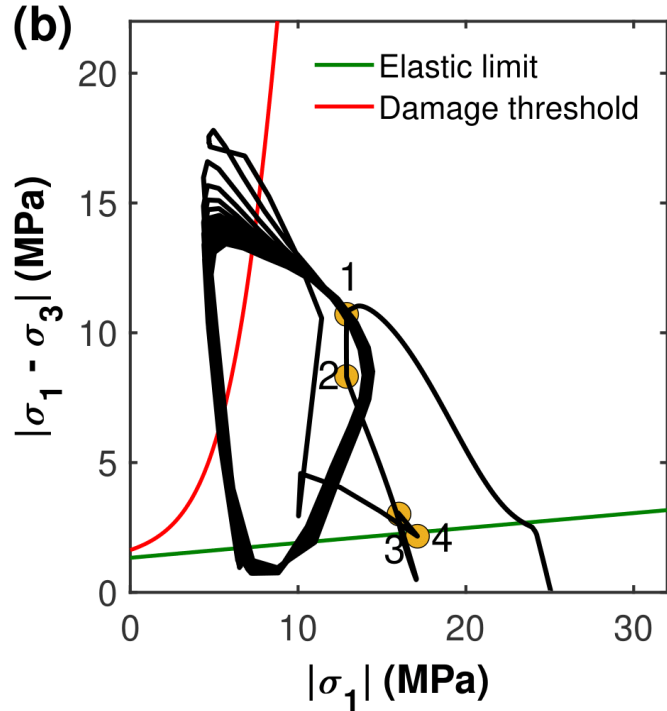
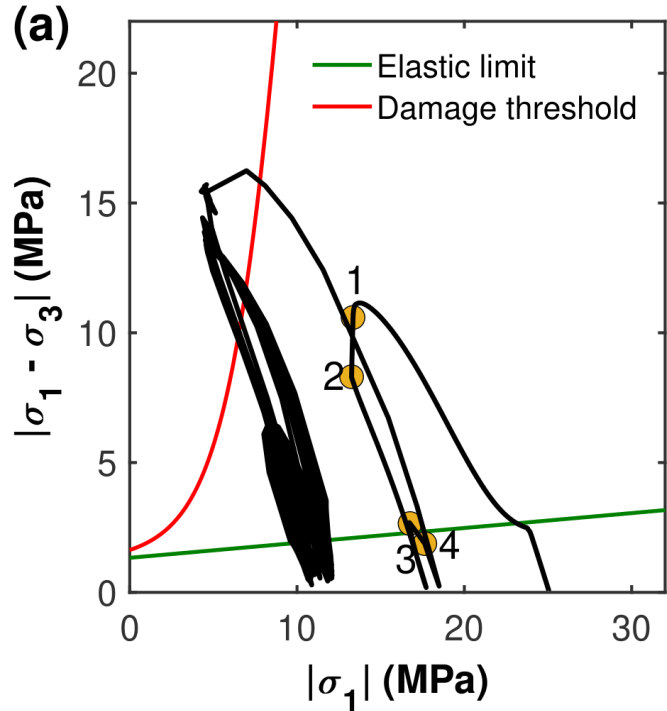


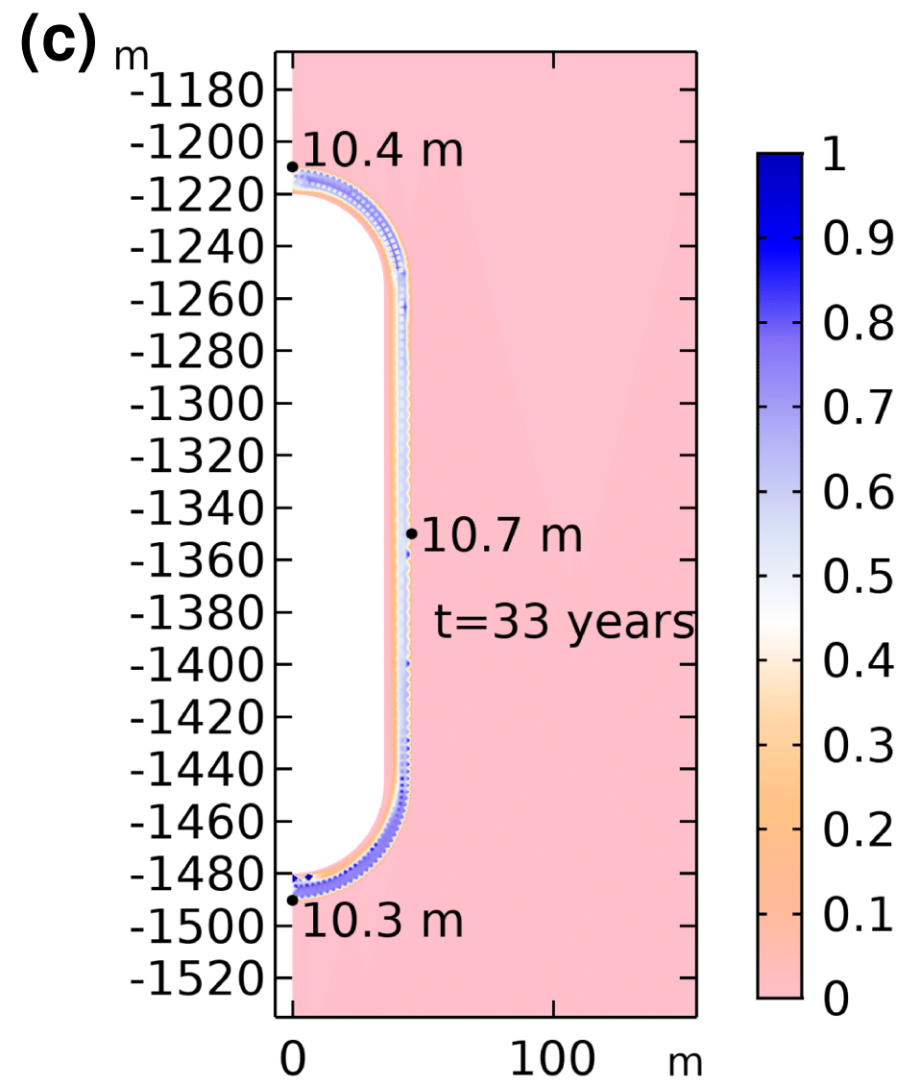
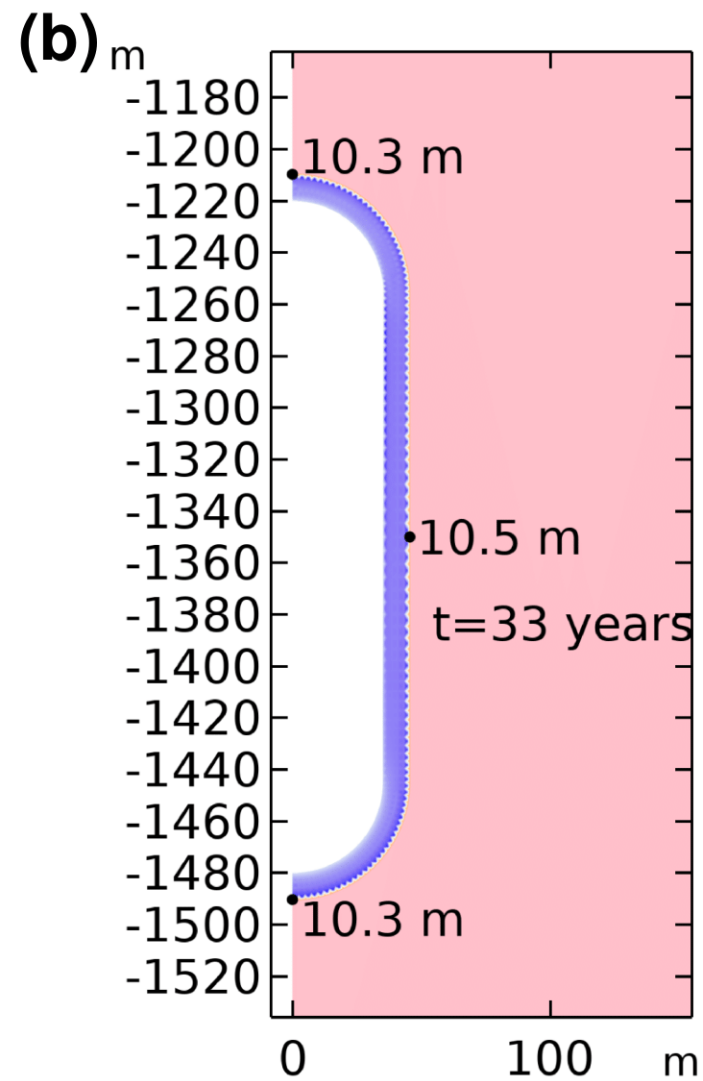
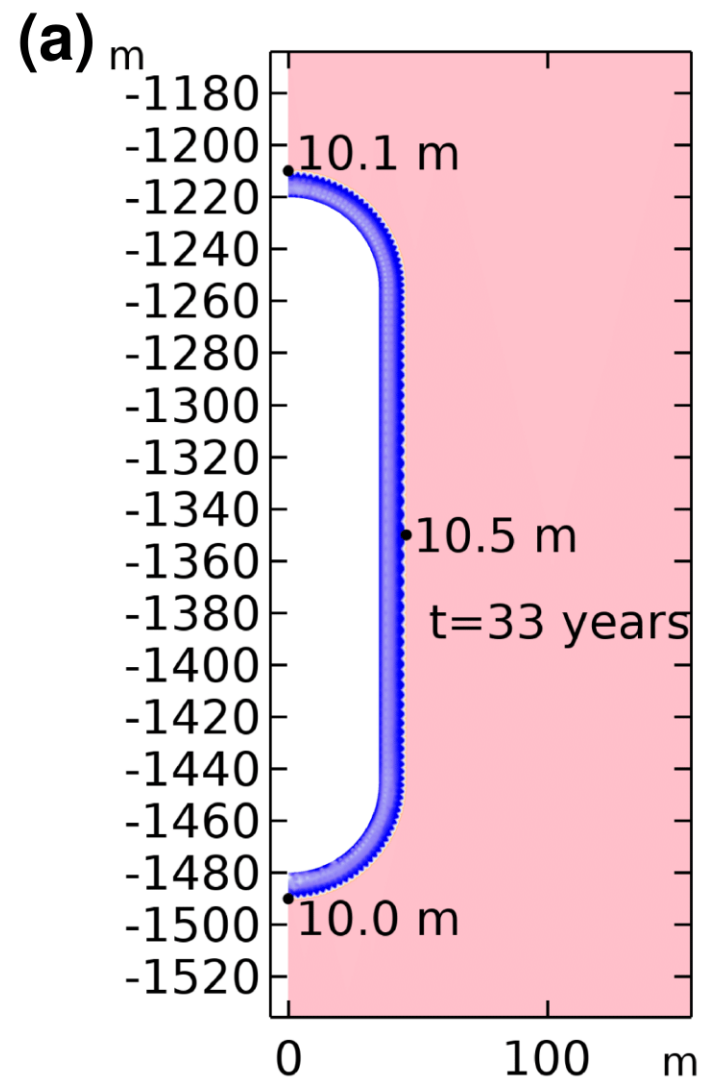
❖ **Damage initialisation:** volumetric dilatation, due to microcracking, lead to a significant increase in the permeability of rock salt

→ Damage initiation is characterised macroscopically by dilatancy

→ The maximum achievable damage is reduced with increasing confining stress, because there is no volumetric dilatancy → ductile behaviour

→ Based on continuum of damage mechanics - CDM, an isotropic damage variable is considered, in first approximation, which modifies the elasticity and strength parameters





Proposed elastoplastic damage model

Full model equations

(compressive stresses are negative and $\sigma_1 \leq \sigma_2 \leq \sigma_3$)

M-C yield surface:

$$f_p = q + pM_p - N_p$$

$$f(p, q, \theta) \quad M_p = \eta^p(\gamma_p) \frac{\sin \phi}{\left(\frac{\cos \theta}{\sqrt{3}} - \frac{1}{3} \sin \theta \sin \phi\right)} \quad ; \quad N_p = \eta^p(\gamma_p) \frac{(1-d)c_i \cos \phi}{\left(\frac{\cos \theta}{\sqrt{3}} - \frac{1}{3} \sin \theta \sin \phi\right)}$$

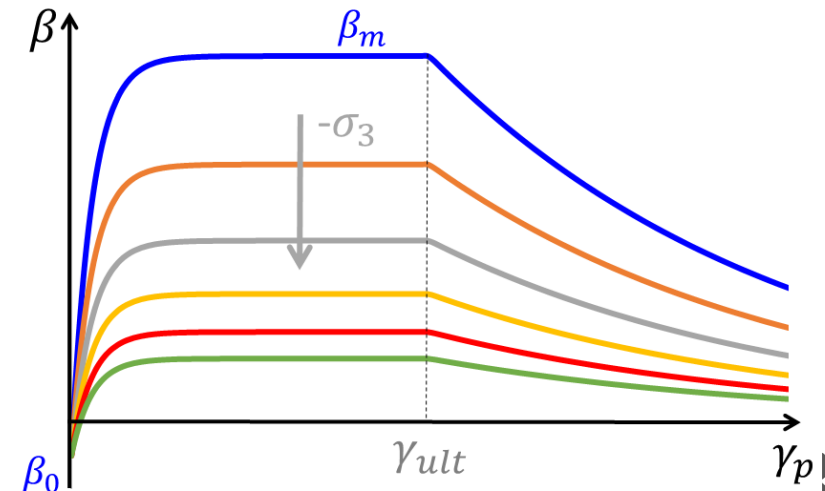
Hardening variable: $\eta^p(\gamma_p) = \left\{ \eta_o^p + \left\{ 1 - \eta_o^p \right\} \frac{\gamma_p}{\alpha^p + \gamma_p} \right\}$

(Chiarelli et al. 2003, Zhou et al. 2011)

Plastic potential: $g = q + \beta(\gamma_p)p$ (Chiarelli et al 2003, Souley et al. 2017)

$$\beta(\gamma_p) = \begin{cases} \beta_m - (\beta_m - \beta_0) \exp(-b_\beta \gamma_p) & ; \quad \gamma_p < \gamma_{ult} \\ \beta_{ult} \exp\left(1 - \frac{\gamma_p}{\gamma_{ult}}\right) & ; \quad \gamma_p \geq \gamma_{ult} \end{cases}$$

$$\beta_m = \beta_{m0} \exp(a_d \sigma_3)$$



Rock salt behaviour in long-term

$$\dot{\varepsilon}_{ij}^c = \left(\frac{1}{\bar{\eta}_k^*} (q^* - \bar{G}_k^* \varepsilon_{tr}) + A_N \left(\frac{q^*}{\sigma_{ref}} \right)^{n_N} \right) \frac{3}{2} \frac{s_{ij}}{q^*}$$

$$\bar{G}_k^* = G_k \exp(k_1 q^*)$$

$$\bar{\eta}_k^* = \eta_k \exp(k_2 q^*)$$

$$q^* = \frac{q}{1 - d_t}$$

q^* : Mises stress undamaged
 q : Mises stress damaged

Transient creep
 (Kelvin-Voigt model)

Parameters: G_k , k_1 , η_k and k_2
 (Heussermann *et al.* 2003)

Steady-state creep
 (Norton law)

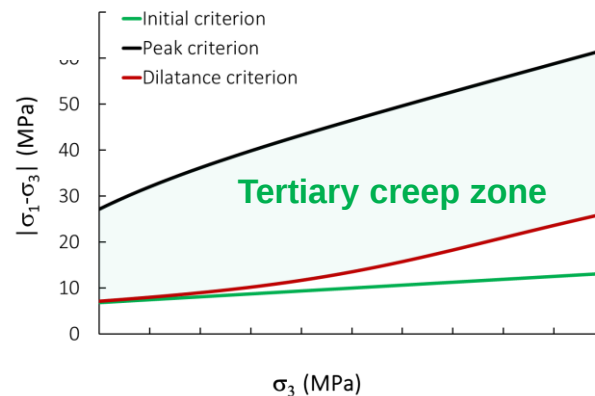
Parameters: A_N and n_N

Tertiary creep
 (Hou *et al.* 2003)

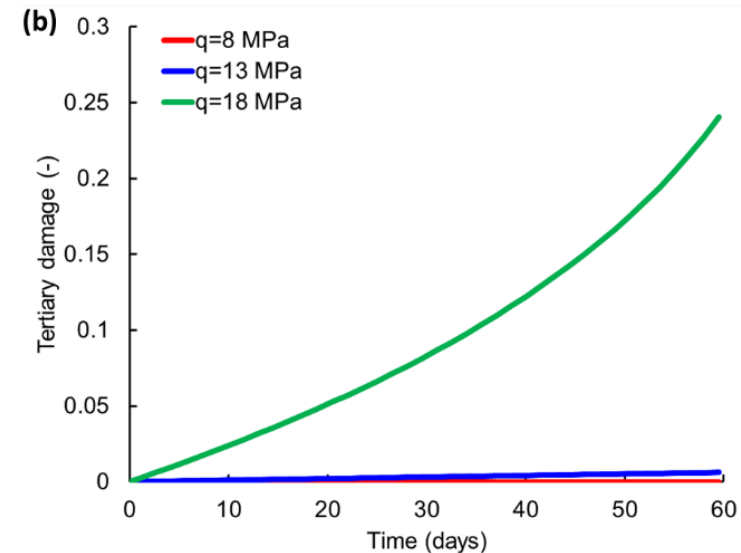
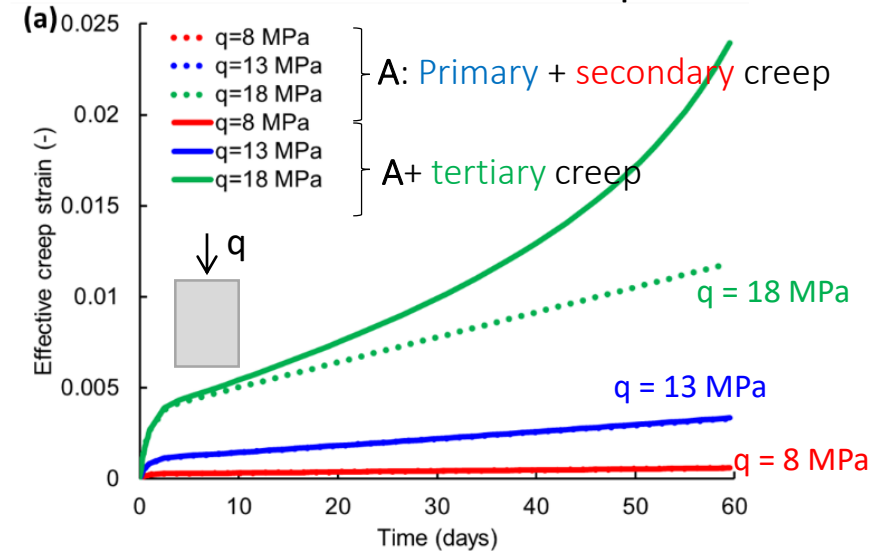
$$\dot{d}_t = A_T \frac{\left(\frac{f_{ds}}{\sigma_{ref}} \right)^{n_T}}{(1 - d_t)^{n_T}}$$

Parameters: A_T
 and n_T

→ Damage is activated and increases if the damage limit
 (dilatance criterion) is exceeded



Numerical uniaxial creep test



Proposed elastoplastic damage model

Full model equations

(compressive stresses are negative and $\sigma_1 \leq \sigma_2 \leq \sigma_3$)

Damage:

$$f_d = d_{max} \{1 - \exp(-b_d(Y_d - Y_0))\} - d \leq 0$$

$$Y_d = \sqrt{\langle \boldsymbol{\varepsilon} \rangle : \langle \boldsymbol{\varepsilon} \rangle}$$

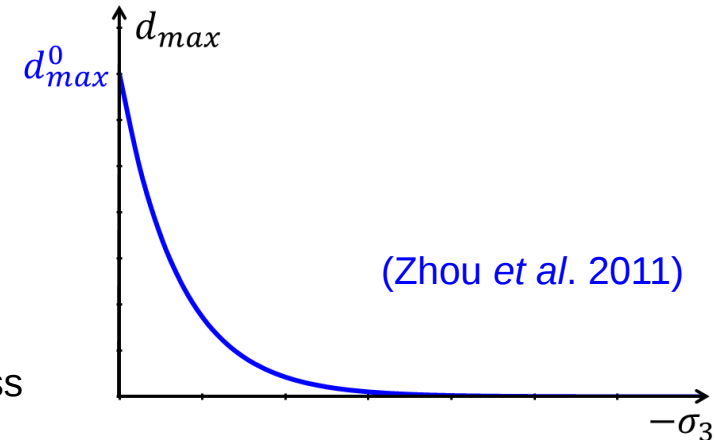
(Mazars
1984)

$$Y_0 = \text{value of } Y_d \text{ at } \beta(\gamma_p^*)=0 \rightarrow f_{ds}$$

$$d_{max} = d_{max}^0 \exp(a_d \sigma_3)$$

a_d : parameter describing the effect of confining stress

d_{max} decreases from its maximum value due to confinement effect



(Zhou et al. 2011)

Dilatance criterion

$$\beta(\gamma_p^*) = 0; \gamma_p^* = -\frac{1}{b_\beta} \ln\left(\frac{\beta_m}{\beta_m - \beta_0}\right) \Rightarrow \eta^p(\gamma_p^*)$$

$$f_{ds} = q + pM_{ds} - N_{ds} \quad M_{ds} = \eta^p(\gamma_p^*) \frac{\sin \phi}{\left(\frac{\cos \theta}{\sqrt{3}} - \frac{1}{3} \sin \theta \sin \phi\right)} \quad ; \quad N_{ds} = \eta^p(\gamma_p^*) \frac{c_i \cos \phi}{\left(\frac{\cos \theta}{\sqrt{3}} - \frac{1}{3} \sin \theta \sin \phi\right)}$$

Parameters:

E, ν

$c_i, \phi, \eta_o^p, \alpha^p, \beta_0, \beta_{m0}, b_\beta, d_{max}^0, b_d, a_d$

Hardening

Plastic
potential

Damage

Thorel 1994 : $c_i, \phi, d_{max}^0, b_d, a_d$

Dragan et al. 2021: $E, \nu, \beta_0, \beta_{m0}, b_\beta$

Numerical verification of triaxial tests

Short-term parameters values used

